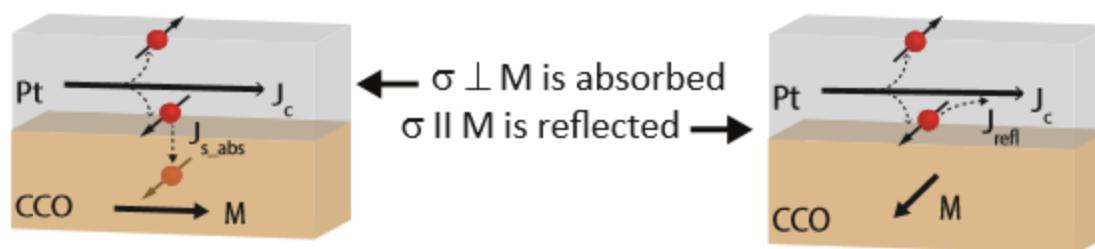


Spin-Hall magnetoresistance

A charge current through Pt creates a spin accumulation by spin-Hall effect (SHE). This spin accumulation can be absorbed or reflected by rotating the magnetization M of CCO [1].

Direction of J_s , J_e and spin polarization (σ), is given by: $J_e \propto (J_s \times \sigma)$



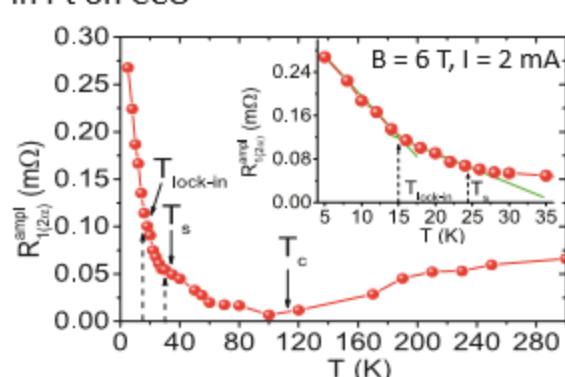
R_1 : 1st harmonic response of transverse Pt resistance

Change in R_1 by rotation of in-plane magnetization M

R_1 : combination of

1. Hall effect (R_H) $\propto I, \sin(\alpha)$
2. SMR (R_{SMR}) $\propto I, \theta_{SH}, \sin(2\alpha)$

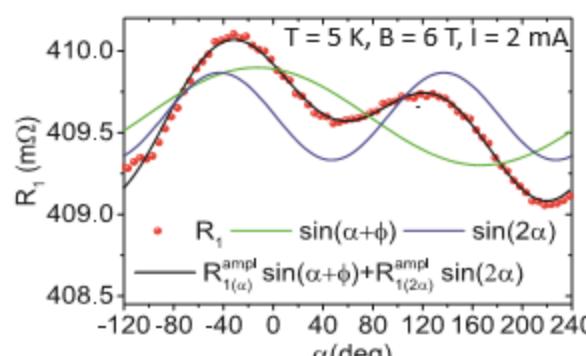
T-dependence of SMR (R_{SMR}) in Pt on CCO



Below T_c in YIG, $R_{SMR} \propto M$. However in CCO, unlike bulk magnetization M,

$$R_{SMR} \propto T^{-1}$$

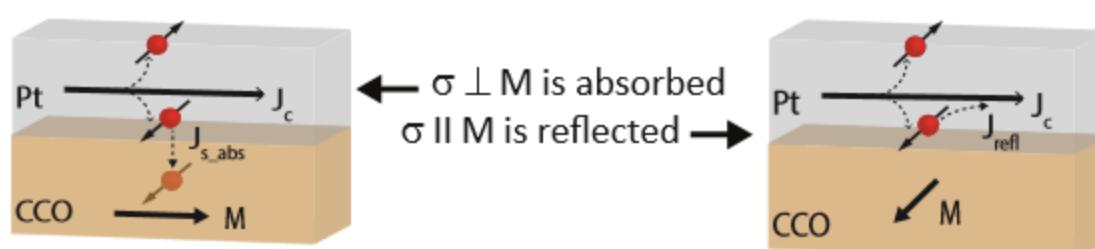
SMR signal in paramagnetic state above T_c can be related to high interface susceptibility.



Spin-Hall magnetoresistance

A charge current through Pt creates a spin accumulation by spin-Hall effect (SHE). This spin accumulation can be absorbed or reflected by rotating the magnetization M of CCO [1].

Direction of J_s , J_e and spin polarization (σ), is given by: $J_e \propto (J_s \times \sigma)$



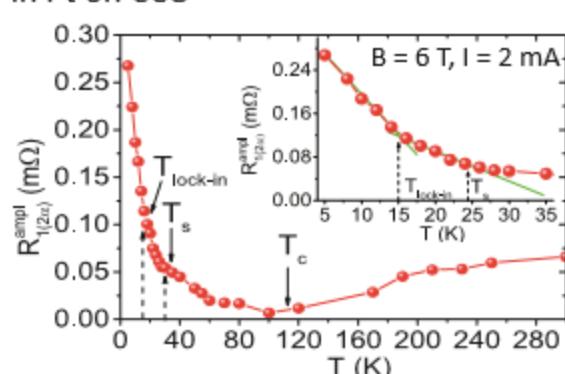
R_1 : 1st harmonic response of transverse Pt resistance

Change in R_1 by rotation of in-plane magnetization M

R_1 : combination of

1. Hall effect (R_H) $\propto I, \sin(\alpha)$
2. SMR (R_{SMR}) $\propto I, \theta_{SH}, \sin(2\alpha)$

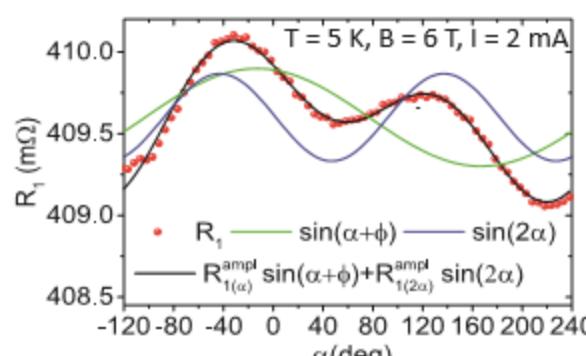
T-dependence of SMR (R_{SMR}) in Pt on CCO



Below T_c in YIG, $R_{SMR} \propto M$. However in CCO, unlike bulk magnetization M,

$$R_{SMR} \propto T^{-1}$$

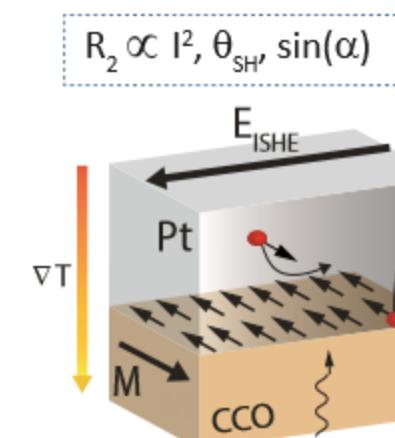
SMR signal in paramagnetic state above T_c can be related to high interface susceptibility.



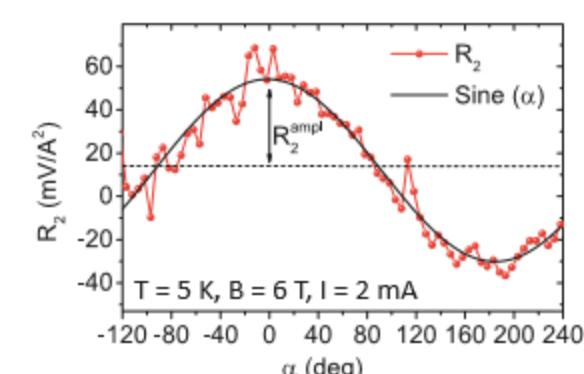
Spin-Seebeck effect

A temperature gradient due to current induced heating over Pt/CCO interface, results in a spin current pumped into the Pt layer, electrically detected by ISHE [3].

R_2 : 2nd harmonic response of transverse Pt resistance measurement

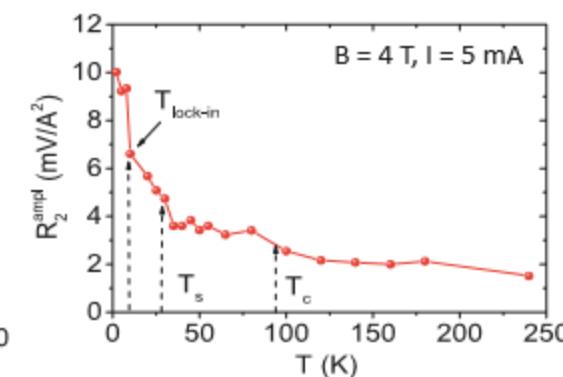
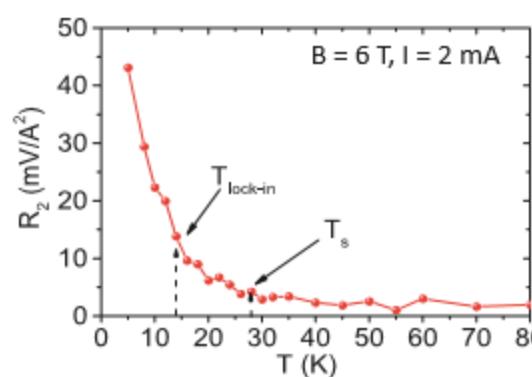


Change in R_2 by rotation of in-plane magnetization M as expected for SSE



T-dependence of SSE (R_2) in Pt on CCO

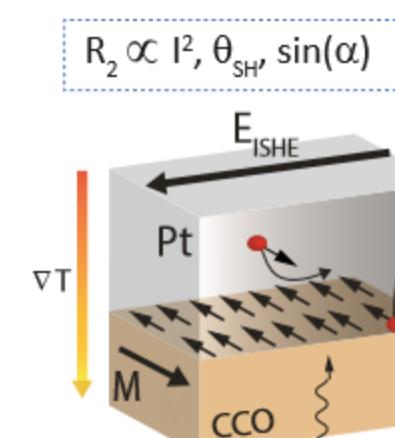
Distinct slope change of SSE (R_2) signal below spin-spiral (T_s) and spin lock-in ($T_{lock-in}$) transition temperatures.



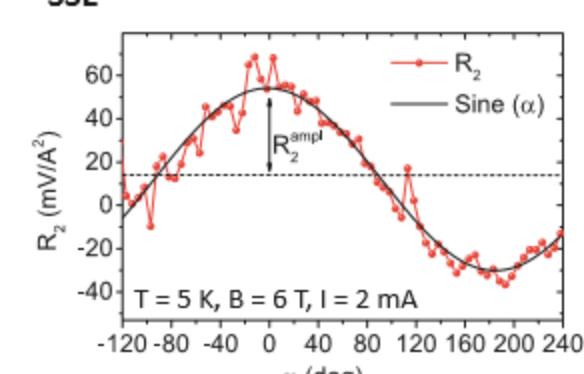
Spin-Seebeck effect

A temperature gradient due to current induced heating over Pt/CCO interface, results in a spin current pumped into the Pt layer, electrically detected by ISHE [3].

R_2 : 2nd harmonic response of transverse Pt resistance measurement

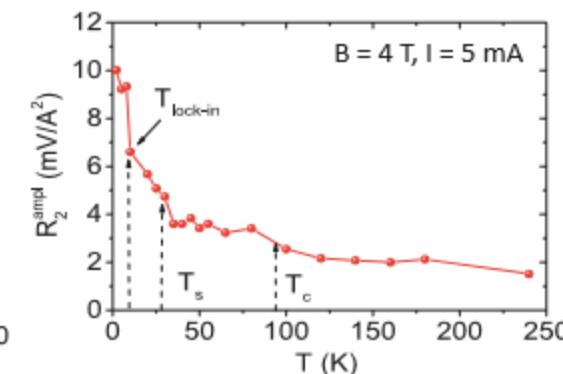
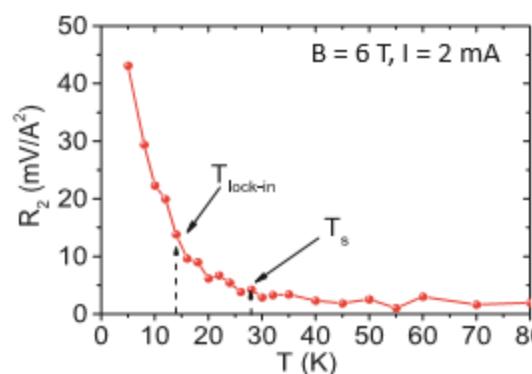


Change in R_2 by rotation of in-plane magnetization M as expected for SSE



T-dependence of SSE (R_2) in Pt on CCO

Distinct slope change of SSE (R_2) signal below spin-spiral (T_s) and spin lock-in ($T_{lock-in}$) transition temperatures.



Summary

- The temperature dependence of the SMR and, though less so, the SSE, exposes distinct anomalies at the magnetic phase transitions.
- The SMR is more than one order of magnitude larger at $T < T_{lock-in}$ as compared to the signals around T_c .
- We relate the observed enhancement of the SMR below T_s to contributions from the cycloidal spiral, projected onto the spin accumulation at the Pt/CCO interface.
- Increase in the SSE signal below T_s suggests that the magnons from the complex Cr-sublattice magnetization texture plays a essential role in the observed signal.

References:

[1]. N. Vlietstra et al., *Phys. Rev. B* **87**, 184421 (2013). [2]. A. Aqeel, et al., *J. Appl. Phys.* **116**, 153705 (2014). [3]. N. Vlietstra et al., *Phys. Rev. B* **90**, 174436 (2014).

Acknowledgments: We are very grateful to Jacob Baas, Henk Bonder, Johan Holstein and Martijn de Roosz for their unconditional and constant support. Submitted on arxiv:1507.01352



university of
groningen

faculty of mathematics and
natural sciences

ernikhe institute for
advanced materials

nanoLabNL

nanonextnl
innovating with micro and nanotechnology