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FORMATION OF SMALL VACANCY CLUSTERS IN TUNGSTEN AROUND SILVER AND INDIUM IMPURITIES STUDIED BY PAC AND THDS

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Vacancy clustering at Ag and In impurities in W was investigated combining Thermal Helium Desorption Spectrometry (THDS) and Perturbed Angular Correlation (PAC) measurements. The results of these experiments are compared with each other and with those obtained from Monte Carlo calculations.

1. Results of PAC measurements

Almost perfectly (100) oriented polycrystalline foils of W were implanted at room temperature with 25 keV and 110 keV ¹¹¹Ag and ¹¹³In ions, at a dose of 2×10¹³/cm² or less. The samples were annealed for 15 min in vacuum at step-wise increasing temperatures, and studied by PAC after each step. From these measurements information was obtained about the symmetry of each lattice defect trapped at the impurity atom, and the number of defects as a function of annealing temperature. For details of the method, see ref. [1] and references quoted therein.

Apart from substitutional impurities, i.e. atoms with no defects within 1 lattice unit (LU), 3 different impurity-defect configurations were distinguished. The corresponding fractions are shown in fig. 1. For ¹¹³InW the results agree with those obtained by Pütz et al. [2].

<table>
<thead>
<tr>
<th>Hyperfine Parameters</th>
<th>Annealing Temperature (PAC)</th>
<th>Assigned Configuration</th>
</tr>
</thead>
<tbody>
<tr>
<td>ω₀ = 0 Mrad/s</td>
<td>400 800 1200 K</td>
<td>¹¹¹AgW</td>
</tr>
<tr>
<td></td>
<td></td>
<td>IV₁</td>
</tr>
<tr>
<td>ω₀ = 133 Mrad/s, η = 0</td>
<td>400 800 1200 K</td>
<td>¹¹³InW</td>
</tr>
<tr>
<td>&lt;&lt;111&gt;</td>
<td></td>
<td>IV₂</td>
</tr>
<tr>
<td>ω₀ = 301 Mrad/s, η = 10</td>
<td>400 800 1200 K</td>
<td></td>
</tr>
<tr>
<td>&lt;&lt;110&gt;</td>
<td></td>
<td>IV₃</td>
</tr>
<tr>
<td>ω₀ = 245 Mrad/s, η = 0</td>
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<td></td>
</tr>
<tr>
<td>&lt;&lt;100&gt;</td>
<td></td>
<td>IV₅</td>
</tr>
</tbody>
</table>

Fig. 1. Annealing behaviour of defects in W, observed in PAC measurements on samples implanted with ¹¹¹Ag and ¹¹³In, respectively. Corresponding hyperfine parameters are indicated on the left-hand side. Assigned impurity-vacancy clusters are given on the right hand side. Annealing time was 15 min.
To facilitate the comparison between PAC and THDS results, the temperature scale of fig. 1 is renormalized according to the annealing conditions in the THDS experiments. The latter (see sec. 2) clearly show vacancy mobility at about 800 K (stage III). Therefore, the formation and break-up of defects observed by PAC in this temperature region have to be associated with trapping and detrapping of vacancies. According to THDS vacancy-type traps in W persist up to 1400 K after implantation with 20-25 keV Ag or In ions, trap concentrations being twice as large as in self-implanted samples. Therefore, the defect observed by PAC at about 1100 K is likely to be a relatively small vacancy cluster, too.

In fig. 2 we give some vacancy configurations that are consistent with the observed quadrupole frequencies ($\omega_0$), asymmetry parameters ($\eta$), and orientations ($\langle..\rangle$) of the defect symmetry axis. In the case of IV$_5$ relaxation of the impurity atom over 0.3 LU along the $<001>$ direction has to be assumed.

2. Results of THDS measurements

A (100) single crystal of W was implanted with 5 keV and 25 keV Ag and In ions, at doses ranging from $2\times10^{12}$/cm$^2$ to $2\times10^{13}$/cm$^2$. The sample was annealed by increasing the temperature to a value $T_a$ at a rate of 40 K/s, and subsequently injected with 250 eV He ions. Defects bind He atoms with characteristic energies which were determined by monitoring the He release during the linear heating. From the number of He atoms desorbing in a certain peak information was obtained about the trap concentration. A survey of the THDS method is given in ref. [3]. The equipment used for the experiments in this study is described in ref. [4].

Some typical helium desorption spectra for AgW are shown in fig. 3. Different peaks are clearly visible. The peaks A-C correspond to He atoms releasing from substitutional Ag atoms (interstitial-type traps). He atoms are bound to vacancy-type traps with much higher binding energies (>2.5 eV in W). Release of these He atoms gives rise to peaks G-H. In fig. 4 the relative concentrations of the two classes of traps are shown as a function of the annealing temperature.

For the In implanted samples no interstitial-type traps were observed, which implies that substitutional In atoms do not bind He atoms at 300 K. We notice that for the samples implanted with 5 keV (20 keV) Ag ions an appreciable number of interstitial-type traps is present only after annealing at 850 K (1350 K). This seems to contradict the PAC results which indicate a nearly 100% substitutional fraction of Ag and In atoms immediately after implantation.

3. Model calculations

A simulation of the binary collision cascade following implantation was performed using the programme MARLOWE [5]. Radial vacancy
distributions around the implanted atom were calculated for 25 keV ions impinging on (100) W at 10° and 0° with respect to the [001] direction. The number of implantation events was limited to 20 in each case. The fraction of the formed Frenkel pairs (FP) that will spontaneously recombine depend strongly on the recombination distance. In the case of 1-3 keV noble gas implantation of Mo, Hou et al. [6] found that for an FP recombination distance of 3.7 LU and a displacement energy of 33 eV calculated vacancy concentrations were in agreement with experimentally observed values. We used parameter values 40 eV and 3.7 LU, respectively, in our calculations for W, the results of which are shown in fig. 5.

A Monte Carlo approach was applied to estimate the fraction of near vacancies migrating towards the implant during early stage III. Only vacancies within a sphere with radius of 4 LU were considered. A jump probability of 1/8 was assumed for each <111> direction. For each non-equivalent lattice point the fate of 2000 vacancies was followed. The vacancy was considered to be trapped once it had jumped a first or second neighbour position with respect to the impurity. The vacancy was considered to be lost for short range capture once it had migrated more than 4 LU from the trap. We found that 56% of the vacancies starting from less than 2 LU were trapped after about 13 jumps, and that 8% of the vacancies starting from 2 to 4 LU were trapped after about 13 jumps. The calculated fractions of In atoms
associated with 0, 1 and 2 extra vacancies in early stage III are 40%, 30% and 15%, comparing fairly well with the experimental PAC values 30%, 20% and 7% (In) and 50%, 15% and 6% (Ag), respectively.

4. Discussion

The most striking difference in the results of the two applied methods is that PAC shows a large substitutional fraction in as-implanted samples, where virtually no substitutional atoms are seen by THDS. Apparently no He atoms are trapped at substitutional Ag atoms because these are effectively shielded by nearby vacancies. In the case of substitutional In atoms, the binding energy is too small to cause trapping at room temperature.

Using the observed annealing temperature $T_a = 550(10)$ K at which vacancy trapping is observed by PAC, the time $\Delta t = 15$ min during which the sample was annealed, the calculated number of vacancy jumps $n = 5-13$ required for short-range capture, and the jump frequency $\nu_o = 4.4 \times 10^{13}/s$ derived from the data given in ref. [7], we calculate the monovacancy migration enthalpy $H_{\Delta V} = kT_a \ln (\nu_o \Delta t/n) = 1.71(5)$ eV.

The defect in W observed by PAC at higher temperatures may be the same as the one suggested by Weidinger et al. [8] for In-implanted Mo. In the latter case the impurity atom was assumed to be surrounded by a nearly tetrahedral cluster of 4 vacancies. However, symmetry and magnitude of the electric field gradient tensor are also consistent with the configuration InVs.

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