POLAR DISK GALAXY FOUND IN WALL BETWEEN VOIDS

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ABSTRACT

We have found an isolated polar disk galaxy in what appears to be a cosmological wall situated between two voids. This void galaxy is unique as its polar disk was discovered serendipitously in an H\textsc{i} survey of SDSS void galaxies, with no optical counterpart to the H\textsc{i} polar disk. Yet the H\textsc{i} mass in the disk is comparable to the stellar mass in the galaxy. This suggests slow accretion of the H\textsc{i} material at a relatively recent time. There is also a hint of a warp in the outer parts of the H\textsc{i} disk. The central, stellar disk appears relatively blue, with faint near-UV emission, and is oriented (roughly) parallel to the surrounding wall, implying gas accretion from the voids. The considerable gas mass and apparent lack of stars in the polar disk, coupled with the general underdensity of the environment, supports recent theories of cold flow accretion as an alternate formation mechanism for polar disk galaxies.

Key words: galaxies: evolution – galaxies: formation – galaxies: kinematics and dynamics – galaxies: structure – large-scale structure of universe – radio lines: galaxies

1. INTRODUCTION

Polar ring galaxies are peculiar galaxies encircled by a ring of stars, gas and dust with perpendicular spin (Rubin 1994). These distinct kinematic components are confirmed by observations of their H\textsc{i} distribution (Schechter et al. 1984; van Gorkom et al. 1987; Arnaboldi et al. 1997), which generally exhibit extended emission throughout the polar ring with a central gap. Kinematically, we see that some of these structures are more disklike than ringlike (Iodice et al. 2006), are massive enough to be self-gravitating (Arnaboldi & Sparke 1994), and have the suggestion of spiral arms (Arnaboldi et al. 1997). The central galaxy is generally a rotationally supported S0 or elliptical galaxy and shares a common geometric center with the polar disk, however is relatively gas poor. Polar disk galaxies are relatively rare, with only a few per cent of S0s that have now or have had a polar ring (Whitmore et al. 1990). They are particularly useful and interesting as their perpendicular dynamics allow a three-dimensional probe of the dark matter potential (Schweizer et al. 1983; Sackett et al. 1994; Combes & Arnaboldi 1996).

The formation mechanism for polar disks is not entirely clear, but presumably requires a second formation event to form the perpendicular spin components. Originally, a major merger or close encounter was assumed necessary (Schweizer et al. 1983), however neither theory resolves the problem of the central disk retaining its rotational support in systems with significantly massive polar disks, as with NGC 4650A (Bournaud & Combes 2003; Bournaud et al. 2005; Iodice et al. 2006). The issue is also confused by indeterminate ages of observed polar disks, which are partly dependent on the stability of polar oriented structures (Sparke 2004). More recent theories and simulations suggest cold accretion along cosmological filaments may form and continue to feed these polar disks (Macciò et al. 2006).

We have found a polar disk galaxy in relatively isolated, void-like conditions, situated in a tenuous sheet-like configuration between two voids. It was found as part of an H\textsc{i} survey of 15 void galaxies selected using an early version of the Watershed Void Finder (Platen et al. 2007) and the Sloan Digital Sky Survey (SDSS; Adelman-McCarthy et al. 2008), with each galaxy in a different void. It has no optical polar ring visible in the SDSS imaging but a massive, extended H\textsc{i} polar disk. This suggests cold flow accretion as the most likely formation mechanism, and sheds a new and surprising light on the nature of galaxies populating the most desolate areas in the Universe (see, e.g., Peebles 2001). Observations are described in Section 2, shown in Section 3, and comparisons with theory made in Section 4.

2. OBSERVATIONS

SDSS J102819.24+623502.6 is part of a pilot for a larger study of galaxy formation in underdense regions of the SDSS redshift survey (K. Stanonik et al. 2009, in preparation). Our lonely galaxy is situated at a R.A. and decl. (J2000) of 10:28:19.02 +62:35:02.0 and a redshift of \( z = 0.0178446 \). It is fairly blue, with a g-band magnitude of 17.67 and r-band magnitude of 17.41 as reported in the sixth data release from the SDSS. Note that these images reach to magnitudes of \( g \sim 23 \), as the drift-scan mode employed by the SDSS results in an effective exposure time in each band of 54.1 s (York et al. 2000). A distance of 76 Mpc is calculated assuming \( H_0 = 70 \) km s\(^{-1}\) Mpc\(^{-1}\), which gives a scale of 1\(\prime\)\(\prime\) = 370 pc.

The 21 cm line of our target was observed for 12 hours at the Westerbork Synthesis Radio Telescope (WSRT) in Maxi- short configuration, which optimizes imaging of very extended sources in a single track observation with shortest baselines of 36, 54, 72, and 90 m. We observed in four intermediate frequencies each with 512 channels, a 10 MHz bandwidth, and two polarizations. The target redshift corresponds to an optical heliocentric velocity of 5302 km s\(^{-1}\), so channel increments of 19.5 kHz give a velocity resolution of 4.27 km s\(^{-1}\).

All calibration was done using standard AIPS procedures. Continuum emission was subtracted from the UV data by linear interpolation from the line-free channels. Image cubes were created with a CLEAN box around the H\textsc{i} emission that cleaned...
down to 2 mJy beam$^{-1}$ (4σ). Images were created with robust weighting using a robustness parameter of 1 to give a balance of high resolution, a beam size of 22$''$1 × 19$''$.5, with low noise, a rms of 0.5 mJy beam$^{-1}$. This corresponds to a 3σ detection limit in column density of 6.2 × 10$^{18}$ cm$^{-2}$ and in mass of 2.0 × 10$^8$ $M_\odot$ for a galaxy with velocity width 100 km s$^{-1}$. All and 1σ moment maps were made, with 1.5σ clipping applied to a mask that is Hanning and Gaussian smoothed over three cells, both in velocity and spatially, respectively.

3. RESULTS

3.1. Galaxy Parameters

The total H$\text{I}$ image and intensity weighted velocity field are shown in Figure 1. We calculate a total H$\text{I}$ mass of $M_{\text{HI}} = 3.0 \pm 0.5 \times 10^9$ $M_\odot$. There is an angle of 90$^\circ$ between the central stellar disk and the polar H$\text{I}$ disk, which warps at a distance of 20 kpc to 77$^\circ$.

A comparison of slices through the data cube in position–velocity along the major axis of the H$\text{I}$ disk compared to the major axis of the central stellar disk (Figure 2) clearly demonstrates that the regular rotation within this misaligned gas disk is completely decoupled from the stellar component. The major axis slice (Figure 2, left panel) confirms the notion from the velocity field that the H$\text{I}$ is a disk and not a ring. The rotation curve increases nonlinearly with increasing radius and appears to flatten out at a radius of about 30$''$, or 11 kpc. The rotation at the last measured point (about 22 kpc) is 100 km s$^{-1}$ (assuming an inclination of 40$^\circ$ derived from a best fit of the velocity field). This implies a dynamical mass of $5 \times 10^{10}$ $M_\odot$ within 22 kpc. The SDSS $r_{90}$ value of 6$''$6 gives a radius of 2.4 kpc, so the extent of the H$\text{I}$ gas is about nine times the extent of the perpendicular stellar disk. The inverse concentration index for SDSS galaxies, taken from the ratio of $r_{50}$ to $r_{90}$, correlates tightly with morphological type (Shimasaku et al. 2001), and in our central, stellar galaxy is consistent with an exponential, rotationally supported disk.

Analysis of the optical spectra available from SDSS allows an approximate measure of the stellar mass in the central galaxy. The parameters are calculated by C.-W. Yip et al. (2009, in
Dobos et al. (2006), and the stellar population model by Bruzual preparation) who utilize the spectrum services provided by Dobos et al. (2006), and the stellar population model by Bruzual & Charlot (2003) to find a mass of \(1.6 \times 10^9 M_{\odot}\) for the part of our galaxy falling within the 3′ fiber. This provides a lower limit for the stellar mass of the entire galaxy. Since this area is roughly half the optical extent of our tiny galaxy, we expect the total stellar mass to be twice that, roughly equal to the \(H_i\) mass in the polar disk. This is similar to the polar disk simulated by Macciò et al. (2006) which had a baryonic mass in the polar ring of only a third the mass of the central S0 galaxy, and NGC 4650A which has twice the mass in gas and stars in its polar disk as in the central host galaxy (Iodice et al. 2002).

Our target was also observed by GALEX as part of its All Sky Survey (AIS), and serendipitously in the field of a guest observer 1619 s near-UV (NUV) exposure. Figure 3 overplots \(H_i\) contours over the deep NUV image. The 19th magnitude UV emission follows the disk of the central optical galaxy, with no detection along the polar disk.

### 3.2. Large Scale Environment

Cosmologically, the galaxy is relatively isolated within a wall between two voids (Figure 4). A cursory look at 18 other polar ring galaxies taken from Whitmore et al. (1990) which are contained within the SDSS redshift survey finds them in average or slightly overdense environments (E. Platen 2008, private communication). This is consistent with a photometric statistical study by Brocca et al. (1997) of the visible environments of 56 polar ring galaxies, which found them to be similar to normal galaxies.

The two voids and their dividing wall were identified with the closely related Watershed Void Finder (WVF; Platen et al. 2007) and Cosmic Spine formalism (Aragón-Calvo et al. 2008). The starting point is the Delaunay tessellation field estimator reconstructed density field inferred from the local SDSS galaxy distribution (Schaap & van de Weygaert 2000). The WVF identifies the voids and traces their spatial outline, the Cosmic Spine formalism recovers the wall forming the boundary between the two voids.

Following its identification, we determine the geometry, shape and orientation of the wall. These are inferred from the principal components of the covariance matrix \(C_{ij} = \text{Var}(r_i r_j)\), computed on the basis of the 34 SDSS galaxies within a distance of 8.5 Mpc around the void galaxy. The wall’s flattening is estimated from the ratio of its smallest eigenvalue, \(e_3\), over the other two eigenvalues \(e_1\) and \(e_2\). The obtained thickness of the wall, \(\sim 1.3\) Mpc, corresponds to a flattening of \(\sim 96\%\).

Figure 4 shows the spatial distribution of the galaxies in and around the wall, with the nearest neighbor 4.7 Mpc away, and the boundary voids roughly 25 Mpc in radius. The size of the upper void is a lower limit as it continues past the edge of the SDSS survey area. These maps are created using galaxy redshifts from SDSS DR6, which is limited to galaxies with \(r < 17.77\) magnitudes and 55″ separation from any neighbors. The polar disk is at an angle of \(\sim 25°\) from the normal of the wall, and the central stellar disk is similarly misaligned \(\sim 25°\) away from the plane of the wall. Roughly, however, the polar disk appears more perpendicular to the wall than aligned.

### 4. DISCUSSION

With our data we can probe in \(H_i\) an area of 650 kpc \(\times\) 650 kpc on the plane of the sky surrounding SDSS J102819.24+623502.6, with a velocity range from 4232 km s\(^{-1}\) to 6196 km s\(^{-1}\). Despite the significant volume probed, no companions are detected above the \(2.0 \times 10^8 M_{\odot}\) limit in \(H_i\). Down to the SDSS limits, the neighborhood surrounding our lonely galaxy is empty within 3.5 Mpc. This environmental isolation makes it unlikely that the substantial \(3 \times 10^9 M_{\odot}\) of neutral hydrogen we detect was accreted through a merger or close encounter, and we propose cold accretion as the most likely mechanism (Binney 1977; Keres et al. 2005; Dekel & Birnboim 2006).

Furthermore, satellite mergers are generally accompanied by a burst of star formation and the addition of existing stars to the system, neither of which are seen. We do see NUV emission in the stellar disk (Figure 3), indicating star formation occurred there within the recent past, however the lack of UV or optical emission in the polar disk is convincing. Perhaps, we are observing this system in the earliest stage of its development,
where the slowly accreted gas has not yet cooled sufficiently to begin forming stars. The measured column density never passes above the $10^{21}$ cm$^{-2}$ surface density threshold proposed by Skillman (1987) even without correcting for inclination, leading us to expect little to no star formation in our extended disk.

Slow accretion has been proposed as a formation mechanism for the optically discovered, H$\alpha$ rich polar disk galaxy NGC 4650A (Iodice et al. 2006). While tidal accretion simulations by Bournaud & Combes (2003) can reproduce massive, gas-rich and star-poor polar rings, Bournaud et al. (2005) note that such a encounters would convert the central galaxy into an elliptical remnant, and not the rotationally supported disk found in NGC 4650A and suggested in our system. Cold flow accretion as a formation mechanism of polar disks has also been “observed” serendipitously in two simulations. Macciò et al. (2006) discovered a polar disk galaxy that formed in their simulation area through the infall of cold gas along a 1 Mpc long filament. Brook et al. (2008) found that their cosmological simulation contained a polar disk galaxy that was stable over timescales of 3–5 Gyr. In both cases, comparisons of the simulated galaxies with observations produce amazing optical and dynamical similarities.

The orientation of the central stellar disk $\sim 25^\circ$ from the wall is surprising. Recent observational and simulation results (Trujillo et al. 2006; Bruno et al. 2007; Aragón-Calvo et al. 2007; Hahn et al. 2007) show that disk galaxies statistically align their spin axis to lie within the wall they are embedded in. Such an alignment would then naturally lend itself to accretion of material within the plane of the wall, forming a polar disk. This, however, is not what we see, and instead the accreted material appears to be falling in out of the voids.

What we may witness reflects the process of a hierarchically evolving void population (Sheth & van de Weygaert 2004). As voids expand and evacuate matter from their interior, they tend to merge with neighboring voids into ever larger underdensities. The heuristic simulation of Dubinski et al. (1993) illustrates how a group of spherical voids, embedded within a larger underdense area, expand until they meet. Once they touch, matter at the sheet-like boundaries of the voids starts to stream along these interstices toward the emerging higher-density boundary surrounding the entire complex of merging voids (see Figure 5). As a result, we observe the thinning and dissolution of the intervoid walls while they get gradually integrated in the emerging larger-scale void and remain visible as its tenuous interior substructure. The position of our galaxy within its wall is reminiscent of this situation, supporting the idea that the polar gas is not falling in from the wall, as that material is instead streaming away, but out of the bounding voids.

We have discovered completely by chance, as in simulations, a key example of polar disk formation in which accretion of cold gas is the most likely formation scenario.

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