Ferroelectricity in the cycloidal spiral magnetic phase of MnWO₄

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We investigate the relationships among magnetic, dielectric, and ferroelectric properties of a frustrated spin system MnWO₄, which undergoes several magnetic phase transitions including a commensurate-incommensurate and a collinear-noncollinear transition. Dielectric and pyroelectric measurements show that the transition into a spiral magnetic ordered phase produces a ferroelectric state. The direction of the electric polarization is perpendicular to the spin rotation axis and the propagation vector of the spiral. These observations agree well with recent theoretical predictions that a cycloidal spiral magnetic ordering can result in electric polarization. In a material where ferroelectricity is induced by magnetic order, we can magnetically tune the ferroelectric transitions by exploiting the difference of the net magnetization between the ferroelectric and neighboring paraelectric phases.

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I. INTRODUCTION

Magnetoelectric multiferroics¹,² are materials that not only show ferroelectric and magnetic order simultaneously, but also display coupling between these properties. The coupling between (ferro)electric and magnetic order is called the magnetoelectric effect and is defined as the induction of magnetization by an electric field or the induction of electric polarization by a magnetic field.³,⁴ Both the magnetoelectric effect as well as multiferroics have been studied for more than four decades. However, the subject has attracted little attention, mainly because of the small value of the coupling between magnetic and ferroelectric order. In most conventional multiferroics the origin of these orders have no relation to each other. The field gained interest in recent years after the discovery of new multiferroics showing a large magnetoelectric effect, such as TbMnO₃,⁵ TbMn₂O₅,⁶,⁷ Ni₃V₂O₈,⁸ and CoCr₂O₄.⁹ A characteristic property of these new multiferroics is the simultaneous occurrence of a ferroelectric phase transition and a magnetic one, which suggests a strong interplay between the magnetic and ferroelectric order.

Recent theoretical studies successfully explain the mechanism of the ferroelectricity in these new multiferroics by means of both phenomenological and microscopic approaches.⁸,¹⁰-¹³ A key factor in the ferroelectricity of these materials lies in their noncollinear spiral magnetic structures with a cycloidal component in which the magnetic structure itself breaks inversion symmetry. The general relation between electric polarization $P$ and magnetization $M$ in systems with spiral magnetic structures has been deduced from symmetry considerations, as described by $P \propto \gamma \epsilon_{ij} \times (\vec{S}_i \times \vec{S}_j)$. Here $\gamma$ is a constant proportional to the spin-orbit coupling and superexchange interactions, $\epsilon_{ij}$ is the unit vector connecting the neighboring $i$ and $j$ sites, and $\vec{S}$ is the magnetic moment. In this equation, $\epsilon_{ij}$ is along the propagation vector of the spiral structure, and $(\vec{S}_i \times \vec{S}_j)$ is parallel to the spin-rotation axis. This equation indicates that a finite electric polarization can appear when the magnetic moments at sites $i$ and $j$ are coupled noncollinearly in a spiral manner, and the spin rotation axis is not parallel to the propagation vector. For example, the cycloidal spiral magnetic structure¹⁴ shown in Fig. 1(a) meets these requirements. In addition, the direction of electric polarization is perpendicular to the spin rotation axis [Fig. 1(a)] and the propagation vector of spiral, and can be reversed by the change in the chirality of the spiral. The collinear sinusoidal structure [$(\vec{S}_i \times \vec{S}_j)=0$; Fig. 1(b)] as well as the screw spiral structure [$\epsilon_{ij}||(\vec{S}_i \times \vec{S}_j)$; Fig. 1(c)] does not induce a net polarization. Microscopically, the origin of the ferroelectricity in spiral spin systems has been reported by Katsura and co-workers.¹⁰ The idea is that spin current (or spin chirality) is induced between noncollinear spins due to Aharonov-Casher effect¹⁵ or an inverted effect of Dzyaloshinskii-Moriya interaction.¹³,¹⁶,¹⁷ These recent studies indicate that most cycloidal spiral magnets could show a multiferroic nature as well as a large magnetoelectric effect, and multiferroics should be seen in a much wider range of materials. In this paper, we examine the ferroelectric and magnetoelectric properties of another spiral magnetic

![FIG. 1. Schematic illustrations of types of magnetic structure with a long wavelength. (a) Cycloidal, (b) sinusoidal, and (c) screw. Geometric configurations of $\epsilon_{ij}$ and $(\vec{S}_i \times \vec{S}_j)$ are also shown for the respective structures.](image-url)
Curie-Weiss constant

axis of magnetization, parallel to the

nated by distorted octahedra of hexagonal-close-packed oxy-

clinic with space group

three magnetic ordered phases; AF1

/moderately frustrated antiferromagnetic

magnetic transitions at

T

appears only below

T

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ments are aligned collinearly along the easy axis of magne-

structure is composed of alternate stacking of the Mn and W

layers along the a axis. Both Mn2+ and W6+ ions are coordi-

ated by distorted octahedra of hexagonal-close-packed oxygen

ions. MnWO4 with magnetic Mn2+ is known to be a

moderately frustrated antiferromagnetic (AF) system. Its

Curie-Weiss constant \( \theta \) is \( \sim -75 \) K, while magnetic order

appears only below \( T_{N3} \approx 13.5 \) K. Thus the frustration parameter

\( f(=\theta/T_{N3}) \) is \( =5 \). Below \( T_{N3} \), there are two more

magnetic transitions at \( T_{N2} \approx 12.5 \) K and \( T_{N1} \approx 6.5-8 \) K, and

three magnetic ordered phases; AF1 (\( T \leq T_{N1} \)), AF2 (\( T_{N1} 

=\leq T \leq T_{N2} \)), and AF3 (\( T_{N2} \leq T \leq T_{N3} \)) phases. In the AF1

phase below \( T_{N1} \), the magnetic structure is commensurate with a

propagation vector \( k_{AF1} = (\pm \frac{1}{4}, \frac{1}{2}, \frac{1}{2}) \). The magnetic

moments are aligned collinearly along the easy axis of magneti-

zation, which is within the \( ac \) plane at an angle of 35°-37°

from the \( a \) axis.\(^{21,22}\) As illustrated in Fig. 2(b), the magnetic

structure of the AF1 phase is characterized by the up-up-
down-down (↑↑↓↓) spin configuration along both the \( a \) and the \( c \) axes, which is often realized in the ground state of the

system with competing magnetic interactions.\(^{23}\) In this paper,

we use an orthogonal setting which differs from the crystal-

lographic monoclinic setting to make the discussion simpler.

We define the orthogonal \( x \), \( y \), and \( z \) axes parallel to the easy

axis of magnetization, parallel to the \( b \) axis, and perpendicular
to both the easy axis and the \( b \) axis, respectively [see Fig.

2(b)]. In the AF2 and AF3 phases, the propagation vector

becomes incommensurate with \( k_{AF2} = (-0.214, \frac{1}{2}, 0.457) \). At

the transition between the AF2 and AF3 phases, the incommen-
surate propagation vector does not change. According to

the model proposed by Lautenschlager and co-workers,\(^{22}\) the

only difference between the two phases is the existence of a

component of the magnetic moments along the \( b(y) \) axis in

the AF2 phase, which does not exist in the AF3 phase. In

more detail, the magnetic structure in the AF2 phase shows an

elliptically modulated noncollinear spiral spin structure,

in which the basal plane of the elliptical spiral contains the

easy axis of magnetization (\( l/x \) axis) and the \( y \) axis. The AF3

phase is a collinear sinusoidal magnetic structure with the

magnetic moments lying along the \( x \) axis. The magnetic

structures of the AF2 and AF3 phases can be categorized in a

cycloidal spiral as in Fig. 1(a) and a sinusoid as in Fig. 1(b),

respectively. Systematic study of the magnetic, dielectric, and

magnetoelastic properties of this material show various

magnetic phases with collinear-noncollinear and incommensurate-commensurate transitions, which reveal that the

spiral magnetic order does induce ferroelectricity.

II. EXPERIMENTAL PROCEDURES

Single crystals of MnWO4 were synthesized by two dif-

ferent crystal growth techniques; flux growth and floating-

zone (FZ) methods. In the flux growth technique, we have
grown single crystals from a Na2WO4 flux, following Ref.

24. A mixture of MnCO3, WO3, and Na2WO4 powder was

placed in a platinum crucible and heated up to 1250 °C. It

was then slowly cooled to 600 °C at a rate of 1 °C/h. The

crystals were removed from the flux by dissolving the flux

material in a NaOH solution. This method resulted in dark

purple crystals with typical dimensions of \( \sim 2 \times 2 \times 2 \) mm.

For the FZ growth technique, first MnWO4 powder was syn-

thesized by calcination of MnCO3 and WO3 at 1050 °C for

60 h in a nitrogen gas flow. The resulting light brown pow-

der was pressed into a rod, which was used in the floating

zone setup. The rod was pulled through at a rate of

5–7 mm/h in an atmosphere of nitrogen using a halogen-
lamp image furnace. The obtained crystals were as large as

\( \sim 6 \) mm \( \times 20 \) mm. The crystals could easily be cleaved per-

pendicular to the \( b(y) \) axis, and had a transparent red-brown

color. The crystals grown by both the techniques were char-

acterized by powder and single-crystal x-ray-diffraction mea-

surements. The powder-diffraction patterns were refined with

GSAS,\(^{25}\) to determine the purity of the crystals. The crystals

were cut, polished, or cleaved to create the desired surfaces

for the measurements of the dielectric constant and pyroelec-

current. Subsequently, two electrodes were formed by

sputtered gold on the widest faces of the crystals. The crys-

tals grown by the gold electrodes were annealed at 400 °C for

2 h. The measurements were done in a Quantum Design

physical properties measurement system (specific heat, di-

electric constant, and pyroelectric current) and a Quantum

Design magnetic properties measurement system (magneti-

zation). We measured the dielectric constant at 1 MHz using

an LCR meter. For the determination of the electric polariza-
III. RESULTS AND DISCUSSION

A. Magnetic and electric properties at zero magnetic field

Powder x-ray-diffraction measurements confirmed that the crystals obtained by both growth techniques adopt the $Pc1/2$ monoclinic structure at room temperature. However, the flux-grown crystals contained a few mole percent of Mn$_2$O$_3$ while no impurity phase was observed in the FZ-grown crystals. In the specific heat data it can be seen how this impurity influences the magnetic phase transitions. Figure 3(a) shows the temperature dependence of specific heat for the crystals grown by the flux and the FZ methods. Three anomalies were observed below 20 K in the specific heat for both crystals. Taking account of results of former studies, the anomalies at 13.5 and 12.5 K are attributed to the phase transitions into the collinear sinusoidal AF3 phase at $T_{N3}$ and the noncollinear elliptical cycloidal spiral AF2 phase at $T_{N2}$. The features of both transitions appear to be little affected by the quality of crystals. However, the shape and temperature of the anomaly ascribed to the transition into the AF1 phase at $T_{N1}$ shows a remarkable difference between the flux and the FZ crystals. For the FZ crystals, a sharp peak is observed at $\sim$7.4 K, which suggests that the transition into the AF1 phase is first order. For the flux grown crystal this transition occurs at lower temperature ($\sim$6.4 K) and becomes broader. The impurities in the flux crystal thus weaken the transition towards the AF1 phase, and stabilize this AF2 phase. This result is consistent with a neutron-diffraction study which revealed that the AF2 phase remains even at $T=1.7$ K in a powder sample containing impurities. Therefore we will present only the data taken for the FZ-grown crystals with higher quality in the following.

Figure 3(b) displays the magnetization ($M$) measured at a magnetic field ($H$) of 0.1 T as a function of temperature. In the $M_x$ ($T$) three anomalies were observed at $T_{N1}=7.4$ K, $T_{N2}=12.5$ K, and $T_{N3}=13.5$ K. However, in $M_y$ ($T$) and $M_z$ ($T$), only two anomalies can be observed at $T_{N1}$ and $T_{N3}$, and no anomaly was observed at $T_{N2}$. In the AF3 phase, the $M_x$ and the $M_z$ exhibit little temperature dependence, whereas $M_y$ gradually decreases with decreasing temperature [see inset of Fig. 3(b)]. This indicates that the magnetic moments develop in the direction along the $x$ axis in the AF3 phase. This is in good agreement with the neutron-diffraction result, which showed that the magnetic moments are aligned along the $x$ axis in AF3. At $T_{N2}$, $M_y$ starts to decrease toward lower temperatures. In the AF2 phase ($T_{N1} \leq T \leq T_{N2}$), only $M_z$ remains nearly constant, whereas $M_x$ and $M_y$ are suppressed upon cooling. Thus the transition from AF3 to AF2 involves only a change in the $y$ component of the magnetization, and the ordered magnetic moments develop along both the $x$ and $y$ axes in AF2. Again, this is consistent with a former neutron result that suggested that an additional component of the moments in the direction along the $y$ axis arises in the AF2 phase. At $T_{N1}$, $M_y$ decreases more rapidly, while the $M_x$ shows a sudden increase. This can be attributed to the disappearance of the $y$ component in $M$. In the AF1 phase, the magnetic moments are again aligned collinearly along the $x$ axis.

In the same temperature range as the magnetization measurements, dielectric and pyroelectric measurements have...
been performed. In both measurements, an electric field was applied along the x, y, and z axes. However, significant anomalies were observed only when an electric field was applied along the y axis. Figures 3(c) and 3(d) show the temperature dependence of the dielectric constant along the y axis ($\varepsilon_y$) and the electric polarization along the y axis ($P_y$) obtained from measurements of the pyroelectric current in the absence of $H$. At $T_{N3}$, no anomaly was observed in both $\varepsilon_y$ and $P_y$. This indicates that the appearance of the AF3 phase with the collinear sinusoidal magnetic structure does not affect the dielectric and ferroelectric properties in MnWO$_4$. Upon cooling, $\varepsilon_y$ exhibits a delta function at $T_{N2}$, which suggests the occurrence of a ferroelectric transition. Indeed, a finite spontaneous polarization develops along the y axis below $T_{N2}$, as shown in Fig. 3(d). With further decreasing temperature, $P_y$ increases gradually into the AF2 phase, which suggests that the transition is second order. At $T_{N1} = 7.4$ K, a stepwise anomaly appears in $\varepsilon_y$, as shown in the inset of Fig. 3(c). $P_y$ discontinuously drops at $T_{N1}$ and disappears in the AF1 phase with the up-up-down-down spin structure.

To demonstrate switching of the spontaneous polarization by an external electric field, i.e., ferroelectricity, we performed the following experiment. First, the sample was poled by applying a positive electric field of $+667$ kV/m above $T_{N3}$, in which the sample was cooled down to 10 K. After poling, the electric field was turned off. Subsequently, a negative (or zero) electric field $E_{rez}$ was applied, which was removed after a while. Then, pyroelectric measurements were performed in the absence of applied electric fields. The results are shown in the inset of Fig. 3(d). When $|E_{rez}|$ is less than $\sim 600$ kV/m, the sign of the pyroelectric current is not changed by the application of the negative electric field. However, the pyroelectric current is reversed when an electric field $E_{rez}$ was applied that was higher than $\sim 600$ kV/m. These results demonstrate that the spontaneous polarization can be reversed by an electric field and the coercive field at 10 K is around $\sim 600$ kV/m. Therefore, it can be concluded that the transition into the AF2 phase leads to ferroelectric order.

The discontinuous change in the electric polarization and the shape of the anomalies in the dielectric constant and the specific heat at $T_{N1}$ suggest that the phase transition from AF1 to AF2 is first order. This can be explained by the discontinuous change in the propagation vector of the magnetic structures. The difference in the magnetic structures between AF2 and AF3 is the component of the magnetic moments in the direction of the y axis, and thus the propagation vector of both phases is the same. Therefore the $T_{N2}$ transition is expected to be of second order, which is consistent with the present observations.

Finite electric polarization, and thus ferroelectric order, appears only in the AF2 phase where the cycloidal spiral magnetic structure has been proposed. In addition, $P_y$ appears in the direction along the y axis, which is perpendicular to both the spin-rotation axis ($\parallel z$ axis) and the propagation vector of the spiral in the proposed spin structure model by Lautenschlager and co-workers. This configuration agrees completely with that predicted by recent theoretical studies, which explain the appearance of ferroelectricity in systems with cycloidal spiral magnetic structures. This result indicates that the ferroelectricity in MnWO$_4$ is attributed to inversion symmetry breaking driven by the cycloidal spiral magnetic order. Conversely, the present result confirms the spin structure model proposed by the former neutron-diffraction study.

The maximum value of $P_y$ in MnWO$_4$ is around 60 ($\mu$C/m$^2$), which is similar to the value of another multiferroic with a cycloidal spiral magnetic structure, Ni$_3$V$_2$O$_8$ ($\sim 100$ ($\mu$C/m$^2$)). Moreover, a similar disappearance of ferroelectricity in the ground state was also observed for Ni$_3$V$_2$O$_8$, though the crystal structures as well as the electronic configurations of magnetic ions of these two compounds are different. From the equation described in the Introduction, we anticipate a larger electric polarization in MnWO$_4$ than in Ni$_3$V$_2$O$_8$. Namely, if we only consider magnetic moments ($S=5/2$ in Mn$^{2+}$ and $S=1/2$ in Ni$^{2+}$) and relative angles between $\vec{S}_1$ and $\vec{S}_2$ [−80° in MnWO$_4$ and −50° in Ni$_3$V$_2$O$_8$], we can expect a larger $P$ in MnWO$_4$. However, the smallness of the orbital momentum of Mn$^{2+}$ ($L=0$) and of the y component of magnetic moments may suppress the magnitude of $P$ in MnWO$_4$. Thus in order to obtain a large electric polarization, we need to take account of all parameters that are included in the equation.

We observed a memory effect below $T_{N1}$, where no substantial electric polarization could be observed. The $P(T)$ data shown in Fig. 3(d) were taken by the following poling procedure. First, the sample was cooled from above $T_{N3}$ while applying a poling electric field ($\sim 250$ kV/m). At 2 K which is considerably below $T_{N1}$, the poling field was removed. Then, the pyroelectric current was measured while heating the sample. If the phase below $T_{N1}$ is truly paraelectric, a single-domain ferroelectric state cannot be realized during the measurement after this poling procedure. However, $P$ can be reversed completely by switching the sign of the poling electric field. In addition, the obtained $P(T)$ curve was the same as that obtained without removing the poling field during the measurement. A probable explanation for the memory effect below $T_{N1}$ is that a small portion of the ferroelectric AF2 phase (e.g., in the form of polar nuclei) still remains below $T_{N1}$ and preserves the polarization direction. Indeed, $P$ did not completely disappear below $T_{N1}$ ($\sim 50$ µC/m$^2$) in a polluted flux-grown crystal where the AF2 phase may remain below $T_{N1}$ to a larger extent than that in an FZ-grown crystal (not shown). Further detailed investigations of the crystallographic and magnetic structures such as diffuse neutron-scattering study are needed to test this hypothesis.

**B. Magnetic field effect on the dielectric and ferroelectric properties**

Figure 4 displays the temperature dependence of $\varepsilon_y$ and $P_y$ in selected magnetic fields applied along the x, y, and z axes. For $H$ along the z axis ($H_z$), the steplike dielectric anomaly at $T_{N1}$ is shifted toward lower temperature with increasing $H_z$ [Fig. 4(a)]. In agreement with the shift of $\varepsilon_y$, the anomaly observed at $T_{N1}$ in $P_y$ also decreases with increasing $H_z$ [Fig. 4(d)]. For $H_{\perp}$ larger than 3 T, both anoma-
lies in $\varepsilon_y$ and $P_y$ disappear down to the lowest temperature. This disappearance can be ascribed to the stabilization of the ferroelectric AF2 phase down to the lowest temperature and the suppression of the AF1 phase by $H_x$. Indeed, $P_y$ at low temperature recovers in $H$ above 3 T, as seen in Fig. 4 of Ref. [24]. In this orientation, a finite electric polarization can be induced by a magnetic field at the lowest temperature region. It should be also noted that the divergent dielectric constant at $T_{N2}$ is suppressed by $H_x$, although the peak position of $\varepsilon_y$ at $T_{N2}$ is slightly shifted $[dT_{N2} (9 T) = 0.06 K]$, as seen in the inset of Fig. 4(a). Corresponding to the suppression of the peak structure in $\varepsilon_y$ at $T_{N2}$, the magnitude of $P_y$ slightly decreases with increasing $H_x$ above 3 T. The effect of $H$ applied along the $y$ axis ($H_y$) is distinctly different from that of $H_x$. As seen in Figs. 4(b) and 4(e), the anomalies in $\varepsilon_y$ and $P_y$ at $T_{N1}$ are shifted toward higher temperature with increasing $H_y$, whereas the anomalies observed at $T_{N2}$ shift toward lower temperature. The polarization data clearly show that the ferroelectric AF2 phase is suppressed by applying $H_y$. However, applying $H$ along the $z$ axis ($H_z$) affects the dielectric and ferroelectric properties less than that along $x$ and $b$, as displayed in Figs. 4(c) and 4(f).

To summarize the effect of magnetic fields on the ferroelectric and magnetoelectric properties, we display the magnetic and electric phase diagrams of MnWO$_4$ for the $H$ along the $x$, $y$, and $z$ axes in Figs. 5(a)–5(c), respectively. The phase boundaries denoted by open triangles, crosses, and open circles were determined from anomalies in the dielectric constant, the pyroelectric current, and the specific heat. The gray areas in the phase diagrams correspond to the phase showing a finite electric polarization (or ferroelectricity). These phase diagrams are very similar to those reported by Ehrenberg and co-workers. We also describe the magnetic phase boundaries obtained by our magnetization measurements as closed circles in the phase diagrams. The comparison of the electric and magnetic phase diagrams clearly demonstrates the strong interplay of the evolution of cycloidal spiral phase and ferroelectric order in MnWO$_4$. 

FIG. 4. (Color online) Temperature profiles of (a)–(c) dielectric constant and (d)–(f) electric polarization along the $y$ axis at selected magnetic fields applied along the $x$, $y$, and $z$ axes for a MnWO$_4$ crystal. The inset of (a) shows overall shapes of $\varepsilon_y$ in the vicinity of $T_{N2}$.

FIG. 5. (Color online) Magnetic and electric phase diagrams of MnWO$_4$ with magnetic fields applied along the (a) $x$, (b) $y$, and (c) $z$ axes. Closed circles, open circles, open triangles, and crosses represent the data obtained by the measurements of magnetization, specific heat, dielectric constant, and pyroelectric current, respectively. Gray regions indicate ferroelectric phases.
By comparing the phase diagrams with the $M(T)$ data in Fig. 3(b), it becomes clear that the shift of the phase boundaries by $H$ are closely related to the difference of the net magnetization between neighboring phases. When $H$ is applied along the $x$ axis [Fig. 5(a)], the paraelectric AF1 phase rapidly vanishes by a fairly low $H$ ($\sim 2$ T). This is related to a large jump of $M_x$ from AF1 to AF2 [see Fig. 3(b)], which indicates that the application of $H_x$ readily stabilizes the AF2 phase. By contrast, the high-temperature ferroelectric transition at $T_{N2}$ is independent of the application of $H_x$ up to 9 T. This feature is consistent with the magnetization data in which $M_x$ is not affected by the transition at $T_{N2}$, as shown in the inset of Fig. 3(b). For $H$ applied along the $y$ axis, the ferroelectric AF2 phase shrinks while the AF1 and AF3 phases expand with increasing $H_y$. This variation also agrees well with the evolution of $M_y$ which is lower in AF2 than in AF1 and AF3 [Fig. 3(b)]. As shown in Fig. 5(c), applying $H$ along the $z$ axis little affects the transition at $T_{N2}$ but shifts $T_{N1}$ slightly toward lower temperature. This can be explained also by the magnetization data: $M_z$ at the AF1 phase is slightly smaller than that in the AF2 phase [Fig. 3(b)]. Thus we conclude that the magnetic control of ferroelectricity in MnWO$_4$ can be realized by exploiting the differences in net magnetization between the ferroelectric and neighboring paraelectric phases.

IV. CONCLUSION

We have grown high-quality single crystals of MnWO$_4$ by means of the floating zone method. Using these single crystals, we observed that MnWO$_4$ is ferroelectric only in the AF2 phase for which a cycloidal spiral magnetic structure has been proposed. In addition, the geometrical relation between the proposed magnetic structure and the direction of the electric polarization observed in the current study agrees well with the prediction of recent theoretical studies on the ferroelectricity in cycloidal spiral magnetic systems. Our study strongly suggests that MnWO$_4$ is a new member in the family of the multiferroic systems, in which the cycloidal spiral magnetic order induces a ferroelectric order. In this compound, the ferroelectric phase expands or shrinks by applying magnetic fields. The shift of phase boundaries, induced by magnetic fields, depends on the difference in the net magnetization between the ferroelectric phase and neighboring paraelectric phases. With increasing difference, the phase boundary responds stronger to applied magnetic fields. This result provides further information to design multiferroics which respond to relatively low magnetic fields. Motivated by intensive studies in the past three years, more experimental and theoretical results confirm magnetically induced ferroelectricity. However, the number of materials demonstrated experimentally is still small. Therefore our result of magnetically induced ferroelectricity in MnWO$_4$ has an important contribution in understanding and developing multiferroics research further.

Note added. Recently, the authors have learned of another similar study of the magnetoelectric study in MnWO$_4$ in a preprint server. The study only focuses on the effect of magnetic fields applied along the $b$ axis. The authors have also recently learned of another similar study.

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The anomalies in $P$ were observed at slightly higher temperatures than those in other properties such as $e$. This is probably due to a higher sweeping rate of temperature (3 K/min) for measurements of the pyroelectric current, which leads to larger temperature gradient between the sample and a thermometer. Hence the transition temperature obtained by the pyroelectric measurements is slightly higher than the actual one.
