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The optimal balance between quality and efficiency in proton radiography imaging technique at various proton beam energies: A Monte Carlo study


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Original paper

Proton radiography is a novel imaging modality that allows direct measurement of the proton energy loss in various tissues. Currently, due to the conversion of so-called Hounsfield units from X-ray Computed Tomography (CT) into relative proton stopping powers (RPSP), the uncertainties of RPSP are 3–5% or higher, which need to be minimized down to 1% to make the proton treatment plans more accurate.

In this work, we simulated a proton radiography system, with position-sensitive detectors (PSDs) and a residual energy detector (RED). The simulations were built using Geant4, a Monte Carlo simulation toolkit. A phantom, consisting of several materials was placed between the PSDs of various Water Equivalent Thicknesses (WET), corresponding to an ideal detector, a gaseous detector, silicon and plastic scintillator detectors. The energy loss radiograph and the scattering angle distributions of the protons were studied for proton beam energies of 150 MeV, 190 MeV and 230 MeV. To improve the image quality deteriorated by the multiple Coulomb scattering (MCS), protons with small angles were selected. Two ways of calculating a scattering angle were considered using the proton's direction and position. A scattering angle cut of 8.7 mrad was applied giving an optimal balance between quality and efficiency of the radiographic image. For the three proton beam energies, the number of protons used in image reconstruction with the direction method was half the number of protons kept using the position method.

1. Introduction

The quality of cancer treatment with protons critically depends on the accurate determination of proton stopping powers (PSP) of the tissues traversed by the protons. Proton treatment planning is based on stopping powers derived from CT images leading to systematic uncertainties of 3–5% and even up to 10% in regions containing bone [1–4]. This may result in a lower dose being deposited in parts of the tumor and/or an overdose in healthy tissues.

The use of proton radiography, which gives direct access to PSPs, can reduce the uncertainty in the translation of the X-ray CT image into a map of PSPs.

Several strategies for proton radiography and its 3D extension proton Computed Tomography (pCT) are being explored by various laboratories [5–10]. Measuring the energy loss and scattering angle of an individual proton passing through the patient is the best way to obtain a sufficiently accurate radiograph of the phantom (patient). For these measurements, the initial proton beam energy should be higher than that chosen for patient treatment to ensure that the protons leave the patient’s body. From the entry and exit points and angles of each proton, combined with the deposited energy, the most likely path can be calculated for each proton. This, together with already available CT data, can improve the accuracy of the PSP prediction, and thereby improve the quality of the proton treatment plan.

At KVI-CART we performed proton radiography (pRG) experiments using a novel prototype of a gaseous PSD (Timepix-based Time Projection Chamber, TPC), developed by Nikhef [8,11]. In these experiments, we used a proton beam energy of 150 MeV to...
irradiate a phantom. The experiments were simulated with Geant4. The comparison of both experimental and simulated data was in excellent agreement, which was a motivation for further studies [8].

The challenge in pRG imaging arises from the MCS of the protons traversing different materials causing blurring of the radiography image. Thus, to improve the image quality and identify each material in the phantom, cuts for the proton scattering angle have to be tuned carefully.

In this paper we present the analysis of two different proton scattering angle calculations, based on direction and position of a proton, and its effect on energy loss radiographs and statistics. The analysis was performed for detectors with different WET, these include: an ideal massless detector, a gaseous detector, silicon (Si) and plastic scintillator detectors. The proton energies chosen for the simulations were 150 MeV, 190 MeV and 230 MeV.

2. Material and methods

2.1. Monte Carlo simulations

The Monte Carlo (MC) simulations discussed in this paper were performed with Geant4 version 9.6.p03 [12]. A scheme of the simulation setup implemented in our code is sketched in Fig. 1.

Each of $5.0 \times 10^7$ generated protons passed through the two PSDs before and after the phantom, and it was stopped in the RED. The proton beam field was generated at a distance of 100 cm from the PSD1 with the size of 6 cm $\times$ 10 cm, in x and y direction, respectively (Fig. 1). No divergence of the beam was applied. Simulations were performed for the WET of each of the PSDs, thus mimicking the various detector types that have been used in pRG/pCT systems with tracking detectors. The list of studied PSD detectors with their specifications is shown in Table 1. Different PSDs types, their materials and thicknesses were taken from setups that were used in proton radiography systems described in the literature [10]. The thickness of the gaseous detector was taken based on the system that we previously used in our experiments at KVI-CART [8], the thickness of the Si detector was taken from the PRIMA I experiment [7], and the thickness of the plastic scintillator was taken from the proton radiography setup used at the Paul Scherrer Institute [13]. The RED in all simulations was filled with BaF2, similarly to what was used during our experiments at KVI-CART.

The proton beam energies used in the simulations are $E_{\text{beam}} = 150$ MeV, 190 MeV and 230 MeV. The energy loss of a proton in the phantom, $\Delta E$, was calculated as a difference between the known proton beam energy, $E_{\text{beam}}$, and the residual energy, $E_{\text{RED}}$, detected in RED:

$$\Delta E = E_{\text{beam}} - E_{\text{RED}}$$

(1)

The reference physics list used in our simulations was QGSP_BIC_EMY, which has been suggested by Cirrone et al. [14] to obtain a high level of accuracy. This physics list has been optimized for hadron radiotherapy simulations.

2.2. Phantom

The phantom is composed of a block of polymethyl methacrylate (PMMA), in which six cylinders of various sizes and materials, including tissue surrogates, were inserted (Fig. 2, left).

Three of the cylinders contain only one tissue surrogate material, such as: liver (In1), breast (In3) and lung (In4) and the other three inserts simulate an air cavity (In2), a lung tumour (In5) and dental fillings (In6). The diameters of the implemented cylinders are 8, 10, 12, 15, and 20 mm for In3, In2, In6, In1, and In4 and In5, respectively. Additionally, the topmost cylinder (In6) is inserted at an angle of 60 degrees with respect to the z-axis (beam direction).

<table>
<thead>
<tr>
<th>PSD detector type</th>
<th>Number of PSD detectors</th>
<th>Material</th>
<th>Physical density (g/cm$^3$)</th>
<th>Material thickness (mm)</th>
<th>WET (mm)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Ideal</td>
<td>1</td>
<td>Air</td>
<td>0.0012</td>
<td>0.001</td>
<td>–</td>
</tr>
<tr>
<td>Gaseous (TPC)[8]</td>
<td>1</td>
<td>Isobutene C$<em>4$H$</em>{10}$</td>
<td>0.0025</td>
<td>30.0</td>
<td>0.394</td>
</tr>
<tr>
<td>Silicon[7]</td>
<td>2</td>
<td>Silicon</td>
<td>2.330</td>
<td>0.4</td>
<td>0.752</td>
</tr>
</tbody>
</table>
2.3. Calculations of proton scattering angles

In proton radiography the images are blurred because of MCS. Thus, it is important to find the best way to reduce the impact of MCS on the extracted proton energy loss radiographic image to minimize the blurring and, consequently, to improve the accuracy of the energy loss map. The scattering angle of a proton was calculated using both the direction and the positions of a proton before and after traversing the phantom, according to the following formulas:

$$\phi_{12}^\delta (\text{rad}) = \cos^{-1} \frac{\vec{p}_1 \cdot \vec{p}_2}{|\vec{p}_1||\vec{p}_2|} \quad (2)$$

where: $\vec{p}_1$ ($p_{x1}, p_{y1}, p_{z1}$) and $\vec{p}_2$ ($p_{x2}, p_{y2}, p_{z2}$) are the proton directions defined with momentum vectors at PSD1 and PSD2, respectively (Fig. 1).

$$\phi_{12}^{\text{pos}} (\text{rad}) = \tan^{-1} \sqrt{\frac{(x_2 - x_1)^2 + (y_2 - y_1)^2}{(z_2 - z_1)}} \quad (3)$$

where: $x_1, y_1, z_1$ and $x_2, y_2, z_2$ represent the proton positions at PSD1 and PSD2, respectively. This represents a “zero order” approximation.

3. Results and discussion

Fig. 3 shows the proton scattering angle distributions determined from the proton direction, $\phi_{12}^\delta$ (Fig. 3a) and the proton position, $\phi_{12}^{\text{pos}}$ (Fig. 3b), respectively, according to (2) and (3). The angular distributions are shown for $E_{\text{beam}} = 150 \text{ MeV}$. The PSD1 and PSD2 were filled with water and the corresponding WET is shown in Table 1. The scattering angle determined from the proton direction (Fig. 3a) shows that a significant number of protons travel in the forward direction, with a minimum scattering, while the distributions based on the proton position (Fig. 3b) is broader.

Various cuts on the scattering angle were applied in order to have a sharper image and to better distinguish the various insertions in the object, as discussed in our previous study [9]. The best trade-off between sharp transitions of various materials in the phantom and the statistics was found for scattering angles smaller than 8.7 mrad. Therefore, this scattering angle was also used in this analysis.

Fig. 4 shows the percentage of protons surviving the 8.7 mrad cut, when proton direction (Fig. 4a) and position (Fig. 4b) are used to make the cut. Using the proton position to calculate scattering angles $\phi_{12}^{\text{pos}}$, more protons, by a factor of two, are accepted to build the radiographic image. Based on the results from [9] we verified...
Fig. 4. Accepted numbers of protons for a proton beam energy of 150 MeV, 190 MeV and 230 MeV after applying an angular cut of 8.7 mrad, using both proton direction (a) and position (b).

Fig. 5. Proton energy loss radiographs of the phantom obtained with 150 MeV protons scattering up to $\theta^\text{pos}_{12} < 8.7$ mrad. Both PSD1 and PSD2 are considered to be (a) ideal detectors (1 $\mu$m thickness, filled with air), (b) WET of a gaseous detector, (c) WET of a silicon detector and (d) WET of a plastic scintillator. The black dashed line in (a) at $X = 1.15$ cm shows the position of the projection in Fig. 6.
that the results for $E_{beam} = 150$ MeV are reproducible. In addition, the number of protons surviving the cut for higher proton beam energies of 190 MeV and 230 MeV were also determined and shown in Fig. 4.

Using the proton position, the number of accepted protons is nearly 50% and increases to 65% and 75% for $E_{beam} = 150$ MeV, 190 MeV and 230 MeV, respectively (Fig. 4b), while using proton direction (Fig. 4a), these numbers correspond to 25%, 31% and 37% for $E_{beam} = 150$ MeV, 190 MeV and 230 MeV. Thus, the statistics shows that it is better to use the proton position method to calculate the proton scattering angle.

Fig. 5 shows radiographic images of the phantom reconstructed based on $D_{12}$ of a proton that passed through all the detectors and that scattered up to 8.7 mrad. The proton energy loss, $\Delta E$, is calculated using (1).

Fig. 5a represents the energy loss radiograph when both PSDs were filled with air (ideal detectors), while Fig. 5b, c and d represent PSDs filled with WET of TPC, Si and plastic scintillators (Table 1). Fig. 5a also indicates with the position of the dashed line, located at $X = 1.15$ cm, for which the projections in Fig. 6a are plotted. The reconstructed energy loss radiographs for PSDs filled with air and WET of TPC are very comparable, Fig. 5a and b, with all inserts visible. When PSDs are filled with Si (Fig. 5c), a part of the radiographic image is blurred when calculated using the proton direction (formula 2). The fluctuation of the images is larger, between 0.95 and 1.1 (15%), with the average value of 1.02 (determined by visual inspection), whereas the fluctuation for plastic scintillator is larger, between 0.97 and 1.31 (34%), with an average value of 1.12.

Similar $\Delta E$ profiles and $\Delta E$ ratios were obtained for the same cut of scattered protons of 8.7 mrad, where the scattering angle $\theta_{12}^{pos}$ was calculated using the proton direction (formula 2). The fluctuations of the profiles and ratios were observed to be larger than in Fig. 5a and b, with the average value of 1.03 for protons with energy of 150 MeV, and can be seen in Table 2. Furthermore, both $\Delta E$ profiles and $\Delta E$ ratios were obtained for higher proton beam energies of 190 MeV and 230 MeV (not shown) and similar trend was observed. The average values of $\Delta E$ ratios are smaller when scattering angle is calculated using the proton position than the proton direction (1.01 for $\theta_{12}^{pos}$ vs. 1.03 for $\theta_{12}^{dir}$, for proton beam energies of 190 MeV and 230 MeV, respectively). All the average values of $\Delta E$ ratio fluctuations are shown in Table 2.

*Fig. 6.* (a) $\Delta E$ profiles through selected material of the phantom obtained with 150 MeV protons scattering with an angle $\theta_{12}^{pos}$ up to 8.7 mrad for different WET (Table 1). (b) $\Delta E$ ratio for 150 MeV protons and $\theta_{12}^{pos}$ up to 8.7 mrad for different WET.

| Table 2 |

Average values of $\Delta E$ ratio fluctuations obtained for a proton beam energy of 150 MeV, 190 MeV and 230 MeV after applying an angular cut of 8.7 mrad, using both proton direction ($\theta_{12}^{dir}$) and position ($\theta_{12}^{pos}$).

<table>
<thead>
<tr>
<th>$\Delta E$ ratio (average)</th>
<th>$\theta_{12}^{pos}$</th>
<th>$\theta_{12}^{dir}$</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>150 MeV</td>
<td>190 MeV</td>
</tr>
<tr>
<td>Gaseous (TPC)</td>
<td>1.03</td>
<td>1.03</td>
</tr>
<tr>
<td>Silicon</td>
<td>1.03</td>
<td>1.03</td>
</tr>
<tr>
<td>Plastic scintillator scintillator</td>
<td>1.12</td>
<td>1.07</td>
</tr>
</tbody>
</table>
4. Conclusions

In this paper we present the results of a study on proton radiography, an imaging technique aiming at a direct (and therefore more accurate) calculation of the PSP of an object under study (patient).

The MCS of protons passing through various materials blurs the energy loss radiographic image. Selecting protons with small scattering angles sharpens the edges between materials in the object, decreasing the blurring. Two ways of calculating the proton scattering angle have been studied, where either the proton direction ($\theta_{\mu}$) or the proton position ($\theta_{\text{pos}}$) was used to calculate the scattering angular distribution. When applying cuts in the radiographic image reconstruction based on $\theta_{\text{pos}}$, the number of “accepted” protons was doubled compared to when $\theta_{\mu}$ was used for all three energies, specifically: 51% vs. 25% with $E_{\text{beam}} = 150$ MeV, 65% vs. 30% with $E_{\text{beam}} = 190$ MeV, and 75% vs. 37% with $E_{\text{beam}} = 230$ MeV.

Plastic scintillators, such as PSD1 and PSD2, seem to be the most disadvantageous in the pRG/pCT imaging system, unless the highest proton beam energy of 230 MeV is used. For a more realistic representation of Si detectors, additional materials should be added to the simulations consisting of, for example, a PCB material to which the Si detector is connected. Protons from the beam pass through the PCB material undergoing MCS and causing further blurring of the radiographic image. Thus, it will make the results for PSD_Si less accurate than for PSD_TPC. Furthermore, the thickness of a single Si detector is usually taken to be at least 300 μm. As PSDs of at least two (or three in some pRG/pCT prototypes) Si detectors are placed before and after the object under study, this will increase the image blurring further. The MCS is element dependent and grows with the atomic number of the material. The most promising PSDs in a pRG/pCT system are expected to be gaseous detectors. The MCS in low-density gas is the smallest amongst the discussed materials giving the least expected blurring of the radiographic image. The readout chip in gaseous PSDs can be located at the surface perpendicular to the beam direction, such that protons will not pass through the chip material, but mostly pass through the gas [8]. More simulations and experiments are required to study the MCS on complex phantoms and real tissues.

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References