A Catalytically Active Membrane Reactor for Fast, Highly Exothermic, Heterogeneous Gas Reactions. A Pilot Plant Study

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Membrane reactors have been frequently studied because of their ability to combine chemical activity and separation properties into one device. Due to their thermal stability and mechanical strength, ceramic membranes are preferred over polymeric ones, but small transmembrane fluxes obstruct a widespread industrial use of a membrane reactor. Consequently, a bench-scale membrane reactor with a tubular, macroporous membrane \( (d_p = 700 \text{ nm}) \) was developed in order to attain increased fluxes. A cooling pipe was concentrically placed inside the tubular membrane to remove heat from the membrane surface, so the present membrane reactor was suitable to conduct exothermic reactions. As a model reaction, the heterogeneous oxidation of carbon monoxide over platinum, with separated feed of carbon monoxide and oxygen, was performed in the present setup. First, the present membrane reactor was characterized by the determination of the transport parameters, structure parameters of the membrane, and the external transfer coefficients. Subsequently fluxes of the reactants and products were measured over a wide range of process conditions. Especially the influence of a transmembrane pressure difference was studied extensively. Furthermore overall conversion of carbon monoxide was measured under various process conditions, and the results were compared with the simulations of a simplified, overall reactor model. From the results of the present investigation, it could be concluded that the application of a pressure difference over the membrane turned out to be a major process control parameter. It increases the product yield and preferentially directs the fluxes toward one side of the membrane. It was shown that even for macroporous catalytic membranes substantial pressure differences are allowed without any slip of unconverted reactants through the membrane. Furthermore, high degrees of conversion were observed in the present setup, and the simulations of the overall reactor model were in reasonable agreement with the experimental data. The overall model contained no adjustable parameters. From this study, the catalytically active, ceramic membrane reactor with separated feed of reactants turned out to be highly flexible and easy to control.

1. Introduction

A membrane reactor is a special type of reactor in which a ceramic or polymeric membrane is an integral part. Both organic (polymeric) and inorganic (ceramic and metal) membranes can be used either as a separation device only or in combination with the catalyst function itself. Polymeric membranes, for instance, are used in chemical reactors in order to create large specific areas (see, e.g., Chen et al., 1992). Owing to their thermal stability, chemical resistance, and mechanical strength, ceramic membranes are often preferred over polymeric ones. Membranes can be used as a catalyst support to combine their separation properties with catalytic activity. Ceramic membranes can be distinguished in porous, such as alumina, and nonporous membranes, such as Pd alloys.

In comparison with porous membranes, nonporous membrane are highly permselective toward one of the components. Therefore, Pd-based membranes have been widely studied in dehydrogenation reactions, because hydrogen is almost exclusively transported through this type of membrane (see, among others, Shu et al., 1991). Due to the thermodynamic limited conversion of dehydrogenation reactions, a Pd membrane can be used for continuous hydrogen withdrawal from the reaction mixture in order to increase the conversion. Itoh (1987) studied the dehydrogenation of cyclohexane to benzene and used a Pd membrane as a hydrogen-permeable wall of a packed catalyst bed. It was experimentally demonstrated that the application of the Pd membrane to separate the hydrogen from the reactants resulted in an increased conversion. Many configurations and applications for Pd-based membranes were studied by Gryaznov (1992) and co-workers (e.g., Gur'yanova et al., 1988). This group also studied the concept of heat removal of an endothermic reaction by an exothermic reaction proceeding at the opposite side of the membrane and vice versa. At one side of the membrane the endothermic dehydrogenation of a hydrocarbon proceeds and the hydrogen that permeates through the membrane reacts at the opposite interface with oxygen to supply the necessary heat. However, nonporous membranes suffer from low transmembrane fluxes, and also additional problems related to metal sintering and mechanical problems obstruct a widespread use of these membranes in chemical reactors. Extensive reviews of Pd membranes in chemical reactors were presented by Shu et al. (1991) and Hsieh (1991).

Porous membranes were also studied in chemical reactors for high-temperature applications, and an overview was presented by Hsieh (1991). Usually, these membranes are made of silica and/or alumina, which have become available with well-defined and narrow pore size distributions (Leenaars et al., 1984, 1985).
Zaspalis et al. (1991a,b) used asymmetric membranes with a microporous, catalytic top layer, for the methanol oxidation reaction (1991a) and the catalytic NO\textsubscript{x} reduction with ammonia (1991b) and discussed several reactor configurations and operation modes. With the exception of the flow-through mode, the transport of reactants through the membrane was normal to the flow direction of the gas. Comparison of the different modes for the methanol oxidation showed that the flow-through mode yielded a higher methanol conversion due to a better utilization of the catalyst by methanol, but a lower selectivity toward formaldehyde (Zaspalis et al., 1991a). From Zaspalis' work it could be concluded that the catalytic membrane reactor provides for a high flexibility and good control.

In comparison with other gases, hydrogen has a high mobility, so the performance of a catalytic membrane reactor (CMR) has been frequently studied for the processing of equilibrium reactions involving hydrogen. Champagnie et al. (1992) presented an experimental study on the dehydrogenation reaction of ethane in a tubular CMR containing a multi-layer membrane with smallest pores of 40 nm—Knudsen transport prevails—and with palladium as a catalyst. Ethane mixtures were supplied to one side of the membrane and the reaction products which permeated the membrane were swept away by an inert gas. In this setup conversions above the equilibrium conversion were obtained, owing to the separation properties of the membrane, which was confirmed by model calculations.

From the aforementioned studies, it can be concluded that a suitable membrane for a CMR should have both a high permeability and a good separation selectivity. To attain increased fluxes Sloot et al. (1990, 1991, 1992) presented a nonpermselective CMR which contained a membrane with a relatively large average pore diameter, $d_p \approx 400$ nm. Moreover, the membrane in their CMR was used to provide a well-defined and controlled catalytic interface between two reactants present at opposite sides of the membrane. This concept of a nonpermselective catalytic membrane reactor with separated feed of reactants (CMRSR) was experimentally tested for the Claus desulfurization reaction and catalytic reduction of NO\textsubscript{x}. Cini and Harold (1991) used this concept in multiphase reaction systems and demonstrated it as a suitable alternative to trickle-bed reactors because it is possible to control the liquid distribution over the catalyst and it avoids the occurrence of hot spots. Sloot et al. (1990, 1992) also presented a flux model to predict the transmembrane fluxes of reactants and products for vapor-phase reactor applications with instantaneous equilibrium reactions, and from their calculations they observed some interesting features of a CMRSR. For a sufficiently high equilibrium constant, a small reaction zone with respect to the membrane thickness appeared to be present inside the membrane. It turned out that small changes in the process conditions of either of the gas phases were compensated for by a shift of the location of this zone, resulting in a (limited) degree of self-controllability (Sloot et al., 1991). Subsequently, it was pointed out that a transmembrane pressure difference could be used to increase the fluxes and simultaneously direct the reaction products preferentially to one side of the membrane. Due to the separation of reactants, a CMRSR has a good controllability and is expected to be a promising alternative for the processing of reactions normally requiring strict stoichiometric feed ratio of reactants.

Catalytic oxidation reactions of hydrocarbons at elevated temperatures are known to proceed fast compared with the transport rates of the components, and, owing to their high exothermicity, are prone to thermal runaway and catalyst destruction. By using a CMRSR as proposed by Sloot et al. (1990, 1991, 1992) for the processing of such reactions, the reaction control may be improved, but, moreover, due to the fact that the transport rates of the reactants and products determine the conversion rate in such a CMRSR (Veldsink et al., 1992, 1993b) thermal runaway is not likely. Furthermore, the separation of the reactants avoids explosive gas mixtures, so the CMRSR can be safely operated at higher concentration levels of the reactants. However, a suitable heat-exchanging device should be present in order to conduct highly exothermic, kinetically fast, heterogeneous gas reactions in a CMRSR.

In the present study, CMRSR for the combustion of hydrocarbons is presented which can be operated at temperatures up to 750 K. The reactor configuration consists of a tubular catalytic membrane with a large average pore diameter, $d_p \approx 1 \mu$m. Concentrically inside the membrane a cooling pipe is placed to remove heat from the membrane surface. A general applicable flux model was developed (Veldsink, 1993), which describes the transport of reactants and products through the membrane accompanied with chemical reactions of arbitrary kinetic rate equations. For a catalytic membrane reactor with separated feed of reactants (CMRSR) slip of unconverted reactants through the membrane substantially reduces the attractive features of this type of reactor. Therefore, the CMRSR possesses a limited operation window, which allows significant simplifications of the transport model (Veldsink et al., 1992). On the basis of this approximate transport model, a more complete reactor model is presented in this study to predict the overall conversion in the present CMRSR as a function of the process conditions.

In the present experimental setup, the heterogeneous combustion of carbon monoxide over platinum was performed as a model reaction. First, fluxes of the reactants and products through the membrane are presented as a function of the process conditions, both in absence or in the presence of a transmembrane pressure difference. Second, overall conversion experiments were carried out, and these results were compared with the predictions from the reactor model.

2. Theory

Fluxes through a catalytic membrane can be calculated from the differential mass balances inside the pores, combining the transport equations, such as the dusty gas model (Mason and Malinauskas, 1983) with the reaction rate equations (see Veldsink, 1993). The steady-state transport rates of components in a porous medium, accompanied with chemical reactions, is calculated from the local conservation of mass for each participating component. For a single component, $i$

\[
0 = -\frac{1}{r} \frac{d(r N_i)}{dr} + R_i \text{ (mol m}_\text{at}^{-3} \text{s}^{-1})
\] (1)

In eq 1, $R_i$ represents the local consumption or production rate according to an arbitrary kinetic rate equation. This kinetic rate expression is usually determined experimentally and is an explicitly known function of the partial pressures of the components and temperature:

\[
R_i = f(T, p_1, ...)
\] (2)

\[R_{\text{inert}} = 0 \]
Figure 1. Typical concentration profiles over the membrane for a fast (line) and an instantaneous (dashes) reaction $A + B \rightarrow$ products.

In eq 1, $N_i$ represents the transport flux through the membrane and should be expressed by the dusty gas model equations (see, e.g., Veldsink, 1993):

$$
\sum_{j=1}^{n} \frac{x_j N_j - x_N N_i}{PD_{ij}^f} = \frac{1}{RT} \frac{dx_j}{dz} + \frac{\sum_{j=1}^{n} x_i B_{ij} P}{PRT} \frac{1}{\mu D_{ik}^f} - \frac{1}{dx_i} \\
\left(\delta_i, \delta_j \right)
$$

with diffusion coefficients

$$D_{ik}^f = \frac{4\sqrt{\pi MT_i}}{SRT}$$

The structure of the porous membrane is taken into account in eq 3 by three parameters, $\epsilon/\tau$, $K_0$, and $B_0$, which generally have to be determined experimentally.

At correctly chosen process conditions, slip of unconverted reactants is avoided in the CMRSR and the reactants will be completely converted inside the membrane. In the absence of a pressure difference over the membrane, typical mole fraction profiles are schematically shown in Figure 1 for a general reaction of $A + B \rightarrow$ products. Figure 1 illustrates that when the CMRSR operates in the desired regime a reaction zone exists which is smaller than the membrane thickness. Increasing the reaction rate causes a contraction of this zone and ultimately results in a reaction plane. In this case the reaction rate is infinitely fast compared with the transport rates of the reactants. It was observed from many model calculations that the transmembrane fluxes became independent of the reaction rate as soon as the reaction zone was smaller than the membrane thickness, and slip of reactants was not observable. Then the fluxes are almost completely determined by the transport rates of the reactants and products, resembling the situation of maximum enhancement in gas-liquid processes of mass transfer accompanied with chemical reactions (Veldsink et al., 1992). Consequently, it is possible to reduce the transport model, presented by eqs 1 and 3:

$$\frac{-\left(P_A - \rho A\right)}{RT} = \frac{K_{0v} \Delta P_A}{RT \ln(r_i/r_j)}$$

If the reaction rate can be assumed instantaneous compared with the transport rates of the reactants, a reaction plane will exist inside the membrane, and reactants $A$ and $B$ (see Figure 1) cannot exist simultaneously. Provided the reaction is sufficiently fast, the calculation of the fluxes for a general reaction scheme $A + v_B B \rightarrow$ products, with reactants $A$ and $B$ present at opposite sides, reduces to the calculation of the physical transport rate of each of the reactants to the reaction plane located at $\delta_r$. Integration of eq 1 in this case leads to analogous flux expressions as eq 5 and the location of the reaction plane, $\delta_r$, can be eliminated from the flux equations by considering that the fluxes of $A$ and $B$ should comply to the reaction stoichiometry. For a diluted, binary gas mixture in the absence of a transmembrane pressure difference, it is possible to derive an explicit expression for the fluxes of the reactants:

$$N_{A,B} = \frac{\left(P_A - \rho A\right)}{RT} = \frac{-K_{0v} \Delta P_A}{RT \ln(r_i/r_j)}$$

In the case when the gas is not diluted and/or a pressure difference over the membrane accounts for an additional convective flow, the fluxes can be calculated from the linearized formulation of the dusty gas model equations, as suggested by Krishna (1987):

$$R_i = \frac{D_A^f}{k_{g,i} \ln(r_i/r_j)} + \frac{1}{\epsilon/\tau D_A^f}$$

and external mass-transfer resistances:

$$R_t = \frac{D_A^f}{k_{g,i} \ln(r_i/r_j)} + \frac{1}{\epsilon/\tau D_A^f}$$

In the case when the gas is not diluted and/or a pressure difference over the membrane accounts for an additional convective flow, the fluxes can be calculated from the linearized formulation of the dusty gas model equations, as suggested by Krishna (1987):

$$n \sum_{j=1}^{n} \frac{x_j N_j}{pd_{ij}^f} = \frac{N_i}{RT} \frac{dx_i}{dz} + \frac{x_i \left(B_{ij} \rho \right)}{PRT} \frac{1}{\mu D_{ik}^f} + \left(\delta_i, \delta_j \right)$$

This approximate model accurately predicted the transport through the membrane as was shown by Veldsink et al. (1993), but unfortunately, an explicit flux equation such as eq 5 cannot be obtained from this equation. However, the unknown mole fractions, $\epsilon/\tau$ as well as the location of the reaction plane itself can be numerically solved from the flux expressions the respective mass balances. Consequently, the basis of the approximate flux model (i.e., eq 6 or 9), an overall reactor model can be developed to predict the overall conversion in a CMRSR as a function of the process conditions.

Prior to the calculations of the fluxes through the membrane, its structural parameters ($B_0, K_0, \epsilon/\tau$) should be determined experimentally from permeation mea-
measurements of a pure gas through the membrane (Mason and Malinauskas, 1983). In this case the transport equation according to the dusty gas model, eq 3, reduces to

\[ N_A = -\frac{1}{RT} \frac{B_0 P}{\mu} + \frac{4}{3} K_0 \frac{8RT}{\pi M_A} \frac{dP}{dz} \tag{10} \]

For a tubular membrane with boundary conditions at \( z = r_i, P = P_i \); \( z = r_o, P = P_o \), and constant flux and physical parameters over the membrane eq 10 results in

\[ N_{Ai} = -\frac{1}{2\mu} \frac{B_0(P_s + P_i)}{2\mu} + \frac{4}{3} K_0 \frac{8RT}{\pi M_A} \frac{(P_s - P_i)}{r_i \ln(r_i/r_o)} \tag{11} \]

According to (11) a plot of the ratio of the permeation flux over the transmembrane pressure difference as a function of the membrane average pressure yields a straight line. From the slope and intercept, the permeability factor, \( B_0 \), and the Knudsen factor, \( K_0 \), can be obtained, respectively. If the pores are supposed as noninterconnected, circular capillaries of radius \( r_p \), the values of \( K_0 \) and \( B_0 \) can be directly calculated (Mason and Malinauskas, 1983) as a function of the pore radius and the tortuosity:

\[ K_0 = \frac{\epsilon \tau}{2\tau}, \quad B_0 = \frac{\epsilon \tau}{8\tau} \tag{12, 13} \]

If this simple structure model holds, the value of the porosity-tortuosity ratio, \( \epsilon \tau \), can be derived from the experimental results of \( B_0 \) and \( K_0 \):

\[ \epsilon \tau = K_0^2 / 2B_0 \tag{14} \]

Since it is not very likely that in porous membranes the pores are perfectly cylindrical, eq 14 must be regarded as an approximation, and the porosity-tortuosity ratio should be determined experimentally and independently from the other parameters (Veldsink, 1993).

Overall conversion of reactants in a CMRSR is the net result of a combination of several transport rates, such as internal pore transport and external gas-to-membrane transport, as well as the degree of back mixing of the gas phases in the CMRSR. To be able to describe the overall conversion in a CMRSR setup, the rates of all transport mechanisms as well as the degree of mixing of the gas phases adjacent to the membrane should be known accurately. In the present CMRSR, the gases are flowing in annular passages at opposite sides of the membrane. Under the present experimental conditions, turbulence was not to be expected so that a Poiseuille velocity profile is present, resulting in a low degree of back-mixing. However, due to the presence of the concentric cool pipe in the tubular catalytic membrane a substantial temperature gradient may exist over the tube-side annular passage, so that natural convection may affect the velocity distribution. On the combined force and natural convective flows in annular passages several investigations have been published (among others, El-Shaarawi and Sarhan, 1981, 1980; Hanzawa et al., 1986a, 1986b) and flow reversal of the fluid along the cold wall was observed under specific process conditions, especially at low \( Re \) (Hanzawa et al., 1986a, 1986b).

A detailed study on the effect of natural convection for the present CMRSR was presented by Veldsink (1993), and it was shown that under the present conditions the velocity distributions were fully developed and almost identical to the Poiseuille profile. In this case the effect of natural convection can be neglected and heat transfer occurs from the hot wall, membrane, to the cooling pipe by a fluid in motion with a fully developed velocity distribution. This is known as the classical Graetz problem, and heat fluxes as well as temperature distributions as a function of the axial distance from the annulus entrance are extensively tabulated for annuli by Lundberg et al. (1963). Consequently, assuming that the Chilton-Colburn analogy of heat and mass transfer (Bird et al., 1966) holds the gas-to-membrane mass transfer coefficient can be calculated from the solutions of the Graetz problem as given by Lundberg et al. (1963).

In the presence of a Poiseuille profile the axial dispersion remains limited. Hence, the overall conversion of the reactants A and B can be calculated from the gas-phase mass balances assuming plug flow. At constant conditions and flow rate the average, mixing cup, the mole fraction of A can be calculated from

\[ \frac{d\chi_A}{dz} = \frac{N_A 2\pi r_p \rho T}{\phi \rho} \tag{15} \]

The approximate model is completed by substitution of either eq 6 or 9 for the fluxes.

3. Experimental Setup

A pilot plant of a catalytic membrane reactor was built which consisted of a tubular, alumina membrane (i.d. = 14 \times 10^{-3} \text{ m}; o.d. = 20 \times 10^{-3} \text{ m}) with an average pore size of \( d_p \approx 700 \text{ nm} \). Two similar membranes were available with respective lengths of, \( L = 0.275 \text{ m} \) and \( L = 0.117 \text{ m} \). The ceramic membrane was connected to a pair of stainless steel (AISI 321) bellows which were subsequently connected to the stainless steel reactor shell. The bellows accomplished a gastight sealing and compensated for the differences in thermal expansion between the ceramic membrane and the metal shell. The connection of the bellows to the membrane was accomplished by Velterop B.V. (Heerhugowaard, The Netherlands). Inside the tubular membrane (tube side) a concentric stainless steel cooling pipe was present (o.d. = 6 \times 10^{-3} \text{ m}) to remove the heat from the membrane surface. Thus, the membrane reactor was divided in two annular gas passages, one enclosed by the membrane and the cooling pipe (tube side, \( b = 4 \times 10^{-3} \text{ m} \) the other by the membrane and the reactor shell (shell side, \( b = 15 \times 10^{-3} \text{ m} \). A schematic drawing of the present membrane reactor is shown in Figure 2. The complete reactor was placed in a temperature controlled oven, which could be used either as a startup device or compensational heating during the experiments. Inlet gas compositions were set and controlled by Brooks’ thermal mass flow controllers and preheated before entering the reactor. The carbon monoxide first passed a fixed bed of copper beads (578 K) to remove any carbonylic compounds. Normally the shell and tube sides of the CMR were operated countercurrently with the tube-side gas passing down flow, although it was easy to realize other configurations. Thermocouples (K-type) were present to monitor the inlet and effluent gas temperatures, membrane temperature, and the temperature of the cooling liquid (Tamson VT4700 oil). Tube and shell side pressures were controlled by back-pressure regulators (Tescam), and the absolute pressures could be read from a digital pressure transducer.
As a model reaction for the characterization of the CMRSR the heterogeneous oxidation of carbon monoxide over platinum was used:

$$\text{CO} + 0.5\text{O}_2 \rightarrow \text{CO}_2 \quad \Delta H_r = -280 \text{ kJ mol}^{-1}$$

4. Results and Discussion

4.1. Hydrodynamics. External, gas-to-membrane mass-transfer coefficients were obtained from the solutions of the Graetz problem for annular passages as given by Lundberg et al. (1963) and by application of the Chilton-Colburn analogy of heat and mass transfer. For the present tubular membrane asymptotic values for fully developed conditions were obtained from

$$\begin{align*}
\text{shell side:} & \quad \frac{k_{g,\text{A}}'2b}{D_A^0} = 4.89; \\
\text{tube side:} & \quad \frac{k_{g,\text{A}}'2b}{D_A^0} = 3.52 \ (16a,b)
\end{align*}$$

As mentioned before, the effects of natural convection and entry length were not included in these asymptotic values, and relations 16 can, therefore, be regarded as a conservative estimation of the external mass-transfer coefficients.

Direct, experimental estimation of the external mass transfer coefficient, $k_g$, from conversion measurements of the heterogeneous carbon monoxide oxidation reaction is possible in the case of premixed reactants at the entrance (Veldsink, 1993). In the present CMR overall conversion of CO was measured as a function of the gas velocity in the tube side at $T = 523$ K, $P = 0.200$ MPa, in order to estimate the gas-to-membrane mass-transfer coefficient. At the tube side, inlet gas mixtures were supplied containing approximately 1.5 vol % CO and 3.2 vol % O2, balance N2, while pure N2 was supplied to the shell side. Under the experimental conditions, conversion of CO ranged from 0.60 to 0.90. The overall mass-flow rates were sufficiently high to assume that natural convection did not affect the flow pattern and a Poiseuille profile was supposed to occur. The experimental results are summarised in Figure 3, plotting the tube-side $k_g$ as a function of the actual volumetric gas-flow rate over the external membrane area, together with the results obtained from eq 16b. From Figure 3 it can be concluded that a good agreement exists between the experimentally estimated $k_g$ and the asymptotic theoretical values as presented in eqs 16.

For the present experiments overall conversions were relatively high, and, as a result, the mass-transfer coefficients that were obtained from these experiments are very sensitive to the occurrence of axial dispersion. Furthermore, the rate of the oxidation of CO may neither be completely instantaneous nor occurring at the external membrane surface only. Owing to these uncertainties and a possible, additional contribution of natural convection, the actual value of $k_g$ may be higher. Therefore, as mentioned before, the external mass transfer coefficient, $k_g$, obtained from eq 16 can be regarded as a conservative estimation.

4.2. Structure Parameters. To determine the structure parameters of the membranes, helium and nitrogen were permeated through the membrane at ambient temperature ($T = 294$ K) and at different pressures. The results were plotted in accordance with eq 11 and Figure 4 shows a typical result. The collective results are summarized in Table 1. The obtained structure parameters were not a function of the direc-

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**Table 1.**

<table>
<thead>
<tr>
<th>No.</th>
<th>$\Delta H_r$ (kJ mol$^{-1}$)</th>
<th>$D_A^0$ (10$^{-10}$ m$^2$s$^{-1}$)</th>
<th>$D_B^0$ (10$^{-10}$ m$^2$s$^{-1}$)</th>
<th>$D_C^0$ (10$^{-10}$ m$^2$s$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>1.52</td>
<td>3.52</td>
<td>2.35</td>
<td>1.03</td>
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<tr>
<td>2</td>
<td>2.52</td>
<td>4.52</td>
<td>3.35</td>
<td>1.53</td>
</tr>
</tbody>
</table>

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**Figure 2.** Schematic drawing of the pilot plant with a tubular membrane.

**Figure 3.** Gas-to-membrane mass-transfer coefficient at the tube side of the membrane plotted as a function of the gas-flow rate. Results were obtained from conversion experiments with premixed feed of oxygen and carbon monoxide at $T = 523$ K and $P = 0.2$ MPa. The line indicates the transfer coefficient from eq 16b.
Permeation measurements of a Pure Gas through the Membrane (Results Summarized for Two Different Membranes)

The permeation measurements of He and N₂ through the short membrane (L = 0.117 m) are plotted according to eq 11. The permeation measurements of CO and O₂ through the membrane have been measured and are plotted as a function of the driving force for non-reaction conditions at 373 K and 0.200 MPa according to eq 6. The transport coefficients were calculated from eq 16, which resulted in values of \( e/\tau = 0.14 \) and \( e/\tau = 0.16 \) from the CO and O₂ diffusion, respectively. In the absence of external mass-transfer limitation the results are \( e/\tau = 0.096 \) and 0.10, respectively. These values are apparently higher than resulting from eq 14, as presented in Table 1, which is explained by the fact that eq 14 holds only for ideal cylindrical pores of uniform size, whereas the actual pore sizes are distributed (see Figure 5). For any distribution it can be shown that always \( x^2 \geq x^2 \) which explains the lower values in Table 1, considering eqs 12 and 13. Furthermore, it appears that the values obtained in the absence of external transport limitations are in closer agreement with the values obtained from the permeation measurements, probably indicating that the actual mass transfer coefficients are higher than calculated from eq 16.

4.3. Transport with Chemical Reaction through a Membrane. 4.3.1. Absence of a Transmembrane Pressure Difference. In a series of experiments, gas mixtures of nitrogen containing small amounts of carbon monoxide passed downflow through the tube side and dilute air countercurrently through the shell side. For these experiments a membrane with \( L = 0.275 \) m was installed in the CMRSR. At a temperature of 523 K, \( P = 0.200 \) MPa, and in the absence of a transmembrane pressure difference, CO fluxes were measured at different pressure ratios of O₂ and CO. The results were plotted in accordance with eq 6 and should yield a straight line from which the effective diffusivities of CO and O₂ could be obtained from its intercept and slope.

Figure 4. Permeation measurements of He and N₂ through the short membrane (\( L = 0.117 \) m) plotted according to eq 11.

Figure 5. Mercury porosigram of the tubular membrane impregnated with Pt/St-Al₂O₃ catalyst. The average pore radius was determined to be 250 nm.

Table 1. Structure Parameters Determined from Permeation Measurements of a Pure Gas through the Membrane (Results Summarized for Two Different Membranes)

<table>
<thead>
<tr>
<th></th>
<th>( L = 0.275 ) m</th>
<th>( L = 0.117 ) m</th>
</tr>
</thead>
<tbody>
<tr>
<td>( 10^{10} \times B_0 ) (m²)</td>
<td>7.721</td>
<td>12.13</td>
</tr>
<tr>
<td>( 10^{9} \times K_0 ) (m)</td>
<td>10.535</td>
<td>12.53</td>
</tr>
<tr>
<td>( K^c_0/(2B_0) = e/\tau )</td>
<td>0.0718</td>
<td>0.0784</td>
</tr>
<tr>
<td>( 10^{9} \times \rho_x ) (m)</td>
<td>293</td>
<td>344</td>
</tr>
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</table>

The \( e/\tau \) ratio was also obtained independently of the other structure parameters from additional diffusion experiments of CO and O₂ through the membrane. At 373 K and 0.200 MPa an inlet gas mixture of CO/N₂ (or O₂/N₂) was supplied to the tube side, while pure nitrogen was supplied to the shell side. Fluxes of CO and O₂ through the membrane have been measured and are plotted in Figure 6 as a function of the driving force in accordance with eq 5. From the results of Figure 7 the overall, transmembrane transfer coefficients of oxygen and carbon monoxide appeared to be

\[
K^c_{\text{CO}} = 1.17 \times 10^{-6}; \quad K^c_{\text{O}_2} = 1.19 \times 10^{-6} \text{ (m²/s⁻¹)}
\]

The porosity-tortuosity ratio was calculated from eq 5. Furthermore, the continuum diffusion coefficients given by Marrero and Mason (1972) and an average value of \( K_0 \) from Table 1 were substituted in eq 5. The transport coefficients were calculated from eq 16, which resulted in values of \( e/\tau = 0.14 \) and \( e/\tau = 0.16 \) from the CO and O₂ diffusion, respectively. In the absence of external mass-transfer limitation the results are \( e/\tau = 0.096 \) and 0.10, respectively. These values are apparently higher than resulting from eq 14, as presented in Table 1, which is explained by the fact that eq 14 holds only for ideal cylindrical pores of uniform size, whereas the actual pore sizes are distributed (see Figure 5). For any distribution it can be shown that always \( x^2 \geq x^2 \) which explains the lower values in Table 1, considering eqs 12 and 13. Furthermore, it appears that the values obtained in the absence of external transport limitations are in closer agreement with the values obtained from the permeation measurements, probably indicating that the actual mass transfer coefficients are higher than calculated from eq 16.

Figure 6. Diffusion of oxygen and carbon monoxide (O) through the membrane at \( T = 373 \) K and \( P = 0.200 \) MPa. Fluxes are plotted as a function of the driving force for non-reaction conditions according to eq 5.

Figure 7. Plot for determining the diffusion coefficient under reaction conditions at 513 K and 0.200 MPa according to eq 6. Regression analysis resulted in values of \( \text{intercept} = 2.60 \times 10^{-2} \) and \( \text{slope} = 1.32 \times 10^{-6} \), respectively.
respectively. Since the conversions were small in these experiments, the partial pressures were assumed to be constant and equal to their logarithmic average. The results of these experiments at low partial pressures of CO are presented in Figure 7. Substitution of the values for the transfer coefficients obtained from eq 16 in eq 6 resulted in effective diffusivities of CO and O₂:

\[ D'_{\text{CO}} = 2.36 \times 10^{-5}, \quad D'_{\text{O}_2} = 1.55 \times 10^{-6} \text{ (m}^2 \text{s}^{-1}) \]  

(17)

From Marrero and Mason (1972) the continuum diffusion coefficients of O₂ and CO under the actual process conditions were obtained as

\[ D^0_{\text{N}_2, \text{CO}} = 2.831 \times 10^{-5}, \quad D^0_{\text{N}_2, \text{O}_2} = 2.747 \times 10^{-6} \text{ (m}^2 \text{s}^{-1}) \]

at \( T = 523 \text{ K} \) and \( P = 0.2 \text{ MPa} \). From eqs 7 and 17 the porosity–tortuosity ratios were determined as \( \epsilon/\tau = 0.14 \) from the CO and a respective value of 0.078 from the O₂ diffusivity. In the absence of mass transfer limitations the respective results for \( \epsilon/\tau \) are 0.11 and 0.063. For the present experiments, heat effects were not observed, and it was experimentally verified that the homogeneous oxidation reactions did not contribute to the overall conversion. The value of \( \epsilon/\tau \) obtained from the effective diffusivity of CO appears to be in good agreement with the results obtained from diffusion experiments in the absence of a chemical reaction, but a substantially lower value resulted from the O₂ diffusivity.

The apparent difference between the diffusivity of O₂ and CO is not readily explained, because the continuum diffusion coefficients of CO and O₂ show the same temperature dependence (Marrero and Mason, 1972), and also the Knudsen diffusivities are nearly identical. Therefore, the effective diffusion coefficients should be approximately equal in both cases. To verify the results, additional experiments of physical transport of oxygen through the membrane, analogous to section 4.2 (see Figure 6), were conducted at \( T = 513 \text{ K} \) and \( P = 0.200 \text{ MPa} \). From these experiments the overall transmembrane transport coefficient (see eq 5) of oxygen appeared to be \( K_{\text{OV}} = 1.72 \times 10^{-6} \text{ (m}^2 \text{s}^{-1}) \), which resulted in \( \epsilon/\tau = 0.14 \) and 0.093 in the presence and absence of external mass-transfer limitations, respectively. Within the experimental accuracy this is sufficiently close to the results from section 4.2. Therefore, the discrepancies between the diffusivities of oxygen and carbon monoxide obtained under reacting conditions from a plot of eq 6 can most likely be attributed to experimental inaccuracies.

From the present flux measurements it can be concluded that the fluxes in the present CMRSR can be accurately predicted by the approximate flux model (eq 6). The values of \( \epsilon/\tau \) obtained from different experiments and conditions are summarized in Table 2, and a reasonable agreement is observed between the diffusion experiments at different process conditions. The lower values resulting from the permeation experiments can be explained from the pore size distribution. Equations 16 give a reasonably accurate, although conservative, estimation of the external mass-transfer coefficient.

### 4.3.2. Presence of a Transmembrane Pressure Difference

The effect of a transmembrane pressure difference on the performance of a CMRSR was extensively studied theoretically by Sloot et al. (1992, 1990) and Veldsink (1993). These studies showed that the presence of a pressure difference over the membrane increases the conversion rate and directs the products toward the low-pressure side of the membrane. To verify these effects for the present reactor setup, experimental fluxes of products and reactants were determined at various pressure differences over the membrane.

First the experiments were carried out at a constant shell-side pressure while the tube-side pressure was increased gradually. Consequently, the CO transport through the membrane was facilitated by the convective flow. The results are presented in Figure 8, in which both the total flux of CO through the membrane and the total outflux of CO₂ (production rate) are presented as a function of the transmembrane pressure difference. In line with the results obtained from the model calculations (Veldsink, 1993) the total CO-transport increases almost proportional with the pressure difference as a result of the contribution of the convective flow to the total transport. Initially, the CO₂ transport increases as a consequence.

The total CO₂ production, however, shows a maximum which can be explained by considering the O₂ transport, which is hindered by occurrence of the convective flow. As a result of the hindered O₂ transport and the facilitated CO transport, the reaction zone (see Figure 1) shifts toward the oxygen side of the membrane, because the transport rates of O₂ and CO should comply to the reaction stoichiometry. Consequently, O₂ has to be transported over a shorter distance to the reaction plane, which apparently compensates the effect of the convective flow, at small pressure differences over the membrane, and results in an increased CO₂ flux.

### Table 2. Collective Results of Porosity–Tortuosity Ratios Obtained from Different Experiments and under Different Process Conditions

<table>
<thead>
<tr>
<th>procedure</th>
<th>( T ) (K)</th>
<th>with ( k_e ) (eq 16a,b)</th>
<th>without ( k_e )</th>
</tr>
</thead>
<tbody>
<tr>
<td>permeation</td>
<td>( \text{N}_2 )</td>
<td>294</td>
<td>0.075</td>
</tr>
<tr>
<td></td>
<td>( \text{He} )</td>
<td>294</td>
<td>0.059</td>
</tr>
<tr>
<td>diffusion</td>
<td>( \text{CO} )</td>
<td>373</td>
<td>0.14</td>
</tr>
<tr>
<td></td>
<td>( \text{O}_2 )</td>
<td>373</td>
<td>0.16</td>
</tr>
<tr>
<td></td>
<td>( \text{O}_2 )</td>
<td>513</td>
<td>0.14</td>
</tr>
<tr>
<td>diffusion + reaction</td>
<td>( \text{CO} )</td>
<td>523</td>
<td>0.078</td>
</tr>
<tr>
<td></td>
<td>( \text{CO} )</td>
<td>523</td>
<td>0.14</td>
</tr>
</tbody>
</table>

\* Values were calculated both in the presence and absence of external mass transfer limitations.
However, at higher pressure differences the increase of the diffusional transport of O₂ cannot compensate for the hindrance of the opposing convective flow anymore, which causes a maximum in the CO₂ production curve. Due to the slip of unconverted reactant, the total CO flux still increases with the pressure difference. Slip of CO, albeit small (<1%), was already observed at ΔP = 60 kPa, which is just before the maximum in the CO₂ production rate curve. As a result of additional conversion at the external membrane surface, the CO₂ production rate still increases before passing through the maximum. Identical results were obtained for an exchange of the CO and oxygen feed. However, Figure 8 clearly shows that application of a transmembrane pressure difference substantially increases the production rate in a CMRSR. Moreover, transmembrane pressure differences up to approximately 50 kPa could be sustained, without slip of reactants, substantially increasing the fluxes of the reactants and the production rate.

A slightly different behavior was observed in the case when the shell-side pressure was gradually increased at a constant tube-side pressure, as shown by Figure 9. Under the presence conditions slip was already observed at ΔP = 40 kPa. Contrary to the previous results no clear maximum for the CO₂ production rate was observed for the present experimental conditions. Similar results were obtained when oxygen and carbon monoxide were exchanged. Probably the difference is caused by a different behavior of the gas volume of the shell and tube side. The tube-side external area of the membrane is smaller than at the shell side and, furthermore, the residence time of the tube side is much smaller. Another possibility might be a nonhomogeneous catalyst distribution over the membrane radius. However, these causes are only speculative, and this point of asymmetric operation is being studied at the moment.

Finally, to illustrate the remarkable effect of a pressure difference over the membrane on the direction of the product fluxes, Figure 10 presents the ratio of the experimental CO₂ fluxes to the tube and shell side of the membrane as a function of the transmembrane pressure difference. From Figure 10 it can be concluded that at the actual conditions application of a pressure difference over the membrane will direct approximately 90% of the product toward the low-pressure side of the membrane before slip of CO occurs. Figure 10 clearly demonstrates that the transmembrane pressure difference is a powerful process variable to control the direction of the reaction products.

4.4. Overall Conversion in a CMRSR. In the present CMRSR the overall conversion of CO that was supplied to the tube side of the membrane in these experiments was measured as a function of the average residence time in this annulus. During these experiments the shell-side conditions remained constant. The actual velocity of the gas was calculated at the average temperature between membrane and cool pipe because of the pronounced temperature difference between the membrane and cool pipe. Figure 11 presents the results of the conversion measurements as a function of the tube-side residence time in the case of a stoichiometric, sub-, or super-stoichiometric inlet partial pressure ratio of CO and oxygen also coincident operation was performed (●).

According to the results presented in Figure 11, a CMRSR allows high degrees of conversion without slip of reactants. Furthermore simulations with the ap-
Conclusions

The present study discussed a newly developed, high-temperature, non-mergeselective, catalytic membrane reactor with separated feed of reactants (CMRSR). The performance of the CMRSR was demonstrated by studying the heterogeneous oxidation of carbon monoxide over supported platinum. Due to the presence of a concentric cooling pipe inside the membrane, heat can be removed from the membrane surface, without affecting the transport rates. Consequently the present CMRSR is suitable for the processing of exothermic reactions, such as combustion reactions.

In the absence of a transmembrane pressure difference the fluxes of CO and O2 were measured at different partial pressure ratios of these reactants. From the results of these experiments it was concluded that, under the present conditions, the overall conversion rates were completely determined by the transport of the reactants. From these experiments the effective diffusivities of CO and O2 could be determined and the results are given by eqs 17. The porosity—tortuosity ratio that resulted from these effective diffusivities appeared to be in good agreement with their values obtained from both permeation measurements and direct diffusion measurements in the absence of a chemical reaction.

It was demonstrated that, even for membrane reactors containing a macroporous membrane, substantial pressure differences over the membrane could be applied. Application of a transmembrane pressure difference favored the conversion rate and directed the reaction products toward the low-pressure side of the membrane. Increasing the pressure difference over the membrane resulted in a maximum in the production rate, which could be explained by considering the transport of the reactant moving against the convective flow.

Finally overall conversion experiments were carried out at different superficial gas velocities at the tube side of the membrane and countercurrent operation of the tube and shell side. Experiments were conducted with different partial pressure ratios of O2 and CO, so the reaction plane inside the membrane was located differently. The results were compared with an approximated reactor model, and the results from this model appeared to be in reasonable agreement with the experimental data. Altering the operation from countercurrent into countercurrent upflow did not affect the overall conversion. The independence of the overall conversion of the operation mode could be explained from the model calculation by considering the location of the reaction plane inside the membrane as a function of the axial membrane distance.

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Nomenclature

\[ A_o \] = external membrane area (m2)
\[ b = (r_1 - r_0) \] = gap width of the annulus (m)
\[ B_o \] = permeability constant (m2)
\[ d = \text{diameter (m)} \]
\[ D_{ij} \] = binary diffusion coefficient between \( i \) and \( j \) (m2 s-1)
\[ D_k \] = Knudsen diffusion coefficient of component \( i \) (m2 s-1)
\[ \text{i.d. = inside diameter (m)} \]
\[ \Delta H_r \] = reaction enthalpy (J mol-1)
\[ k_g \] = gas-to-membrane mass transfer coefficient (m s-1)
\[ K_o \] = Knudsen constant (m)
\[ K_{ov} \] = overall transport coefficient, defined by eq 5 (m2 s-1)

Table 3. Experimental Conditions for the Overall Conversion Experiments Presented in Figure 11*

<table>
<thead>
<tr>
<th>[ 10^{-4} \times \phi_{i}^0 ]</th>
<th>[ x_0^0 ]</th>
<th>[ x_{CO}^0 ]</th>
<th>[ L ]</th>
<th>[ T_m ]</th>
<th>[ P_i ]</th>
<th>[ P_{out} ]</th>
</tr>
</thead>
<tbody>
<tr>
<td>[ (m^3 s^{-1}) ]</td>
<td></td>
<td></td>
<td>(m)</td>
<td>(K)</td>
<td>(MPa)</td>
<td>(MPa)</td>
</tr>
<tr>
<td>1.78</td>
<td>0.010</td>
<td>0.020</td>
<td>0.275</td>
<td>623</td>
<td>0.20</td>
<td>0.20</td>
</tr>
<tr>
<td>1.012</td>
<td>0.109</td>
<td>0.050</td>
<td>0.117</td>
<td>623</td>
<td>0.20</td>
<td>0.20</td>
</tr>
<tr>
<td>1.58</td>
<td>0.013</td>
<td>0.10</td>
<td>0.117</td>
<td>623</td>
<td>0.20</td>
<td>0.20</td>
</tr>
</tbody>
</table>

* The structure parameters of the membrane are presented in Table 1. Conditions are presented for stoichiometric (CASE I), substoichiometric (II), and superstoichiometric partial pressure ratios, respectively.

Figure 12. Location of the reaction plane as a function of the axial coordinate. Results are shown for countercurrent operation (line) and cocurrent operation (dashes). Model parameters: \[ T_m = 623 \text{ K}, P_i = P_o = 0.20 \text{ MPa}, x_{CO} = 0.05, x_0 = 0.025, \phi = \phi_0 = 1 \times 10^{-5} \text{ Nm}^{2} \text{ s}^{-1}, D_{CO} = 1.65 \times 10^{-5} \text{ m}^{2} \text{ s}^{-1}, D_{CO}^N = 2.52 \times 10^{-4} \text{ m}^{2} \text{ s}^{-1}. \]
\[ L = \text{length of the membrane} \]
\[ M = \text{molar mass (kg mol}^{-1}\text{)} \]
\[ N = \text{molar flux (mol m}^{-2}\text{ s}^{-1}\text{)} \]
\[ \text{o.d.} = \text{outside diameter (m)} \]
\[ P = \text{pressure (Pa)} \]
\[ P_j = \text{partial pressure of component } j \text{ (Pa)} \]
\[ R = \text{universal gas constant (J mol}^{-1}\text{ K}^{-1}\text{)} \]
\[ R_s, R_t = \text{dimensionless, external mass-transfer resistance (see eq 8)} \]
\[ r = \text{radius (m)} \]
\[ T = \text{temperature (K)} \]
\[ V = \text{volume (m}^3\text{)} \]
\[ x = \text{mole fraction} \]
\[ \bar{x} = \text{radical averaged mole fraction of A} \]
\[ z = \text{space coordinate (m)} \]

Greek Symbols
\[ \beta = \text{thermal expansion coefficient (K}^{-1}\text{)} \]
\[ \delta_i = \text{location of (hypothetical) reaction plane inside the membrane (m)} \]
\[ \epsilon = \text{porosity} \]
\[ \tau = \text{tortuosity} \]
\[ \zeta = \text{relative conversion} \]
\[ \mu = \text{viscosity (Pa s)} \]
\[ \phi = \text{density (kg m}^{-3}\text{)} \]
\[ \phi_2 = \text{volumetric flow rate (m}^3\text{ s}^{-1}\text{)} \]

Subscripts / Superscripts
\[ A, i, j = \text{with respect to component } A, i, \text{ or } j \]
\[ \text{cool} = \text{cooling pipe} \]
\[ \text{eff} = \text{effective} \]
\[ \text{hps} = \text{high-pressure side} \]
\[ \text{lps} = \text{low-pressure side} \]
\[ \text{m} = \text{average} \]
\[ \text{membrane} \]
\[ p = \text{pore} \]
\[ 0, L = \text{at the beginning and end of the membrane} \]
\[ t = \text{tube side of the membrane} \]
\[ s = \text{shell side of the membrane} \]
\[ 0 = \text{gas phase} \]

Literature Cited


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